

Neutral meson mixing and related observables in the $D_{(s)}$ and $B_{(s)}$ meson systems

Justus Tobias Tsang

for the RBC-UKQCD Collaborations

based on **arXiv:1812.08791**

Wuhan, Lattice2019

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THE UNIVERSITY *of* EDINBURGH



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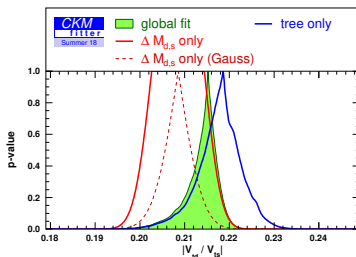
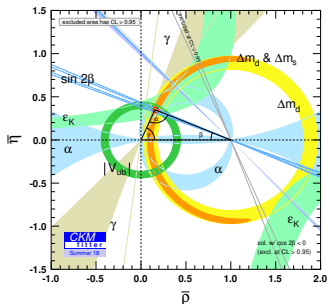
Chris Sachrajda

Outline

- 1 Introduction
- 2 Results for SU(3) breaking ratios (**arXiv:1812.08791**)
- 3 Ongoing Work
- 4 Conclusion and Outlook

Motivation for charm and bottom flavour physics

- Huge experimental efforts: LHC, Belle II, BES III, ...
- Constrain CKM unitarity by combining non-perturbative input with experimental data.
- Test CKM matrix by determining the same CKM matrix element from different processes
- Constrain BSM models
- Address lepton flavour universality (violations?)



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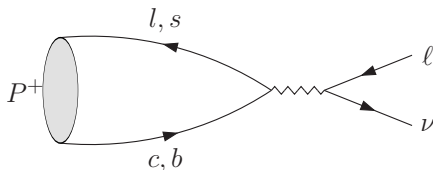
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Related RBC/UKQCD talks:

- Mon 15:40 F. Erben:
"An exploratory study of heavy-light semi-leptonics using distillation"
- Mon 16:50 R. Hill:
"Semi-leptonic B decays with RHQ b quarks"
- Poster O. Witzel:
"Semi-leptonic form factors for exclusive $B_s \rightarrow K\ell\nu$ and $B_s \rightarrow D_s\ell\nu$ decays"

Flavour Physics and CKM: leptonic decay constants

Experiment \approx CKM \times Lattice \times (PT+kinematics)



Leptonic decays: $\Gamma(P \rightarrow l\nu_l) \approx |V_{q_2 q_1}|^2 \times f_P^2 \times \text{known factors}$

where $\mathcal{Z}_A \langle 0 | \bar{c} \gamma_4 \gamma_5 q | D_q(0) \rangle = f_{D_q} m_{D_q}$, $q = d, s$

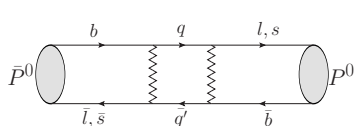
[HFLAV+BESIII] $f_D |V_{cd}| = (45.9 \pm 1.1) \text{ MeV}$, $f_{D_s} |V_{cs}| = (249.1 \pm 3.2) \text{ MeV}$

Computing f_{D_s}/f_D gives access to V_{cs}/V_{cd}

Neutral meson mixing

Neutral mesons oscillate with their antiparticles:

⇒ Difference between mass eigenstates: Δm^{exp} measured to $< 1\%$!



$$\mathcal{A} \propto \left| \sum_{q=u,c,t} \frac{m_q^2}{M_W^2} V_{qb} V_{ql}^* \right|^2 \approx \frac{m_t^4}{M_W^4} |V_{tb} V_{tl}^*|^2$$

$$\Delta m \propto \underbrace{\langle B_{(s)}^0 | \mathcal{H}^{\Delta b=2} | \bar{B}_{(s)}^0 \rangle}_{\text{Short distance}} + \underbrace{\sum_n \frac{\langle B_{(s)}^0 | \mathcal{H}^{\Delta b=1} | n \rangle \langle n | \mathcal{H}^{\Delta b=1} | \bar{B}_{(s)}^0 \rangle}{E_n - M_{B_{(s)}}}}_{\text{Long distance}}$$

$$m_t^2 V_{tb} V_{tl}^* \gg m_c^2 V_{cb} V_{cl}^* \gg m_u^2 V_{ub} V_{ul}^* \Rightarrow \text{Short distance dominated.}$$

Operator Product Expansion

Two scale problem: $\Lambda_{\text{QCD}} \sim 1 \text{ GeV} \ll m_{EW} \sim 100 \text{ GeV}$:

\Rightarrow Factorise via OPE

$$\Delta m \propto \sum_i C_i(\mu) \langle B_{(s)}^0 | \mathcal{O}_i^{\Delta b=2}(\mu) | \bar{B}_{(s)}^0 \rangle$$

- Perturbative model-dependent Wilson coefficients $C_i(\mu)$
- Non-perturbative model-independent matrix elements of $\mathcal{O}_i^{\Delta b=2}(\mu)$
- 5 independent (parity even) operators \mathcal{O}_i .

\Rightarrow SM: $\mathcal{O}_1 = (\bar{b}_a \gamma_\mu (\mathbb{1} - \gamma_5) q_a) (\bar{b}_b \gamma_\mu (\mathbb{1} - \gamma_5) q_b) = \mathcal{O}_{VV+AA}$
+ 4 BSM operators: $\mathcal{O}_2 - \mathcal{O}_5$

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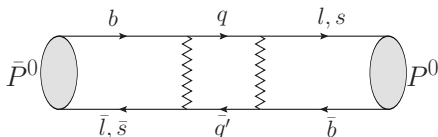
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RBC/UKQCD's $K - \bar{K}$ BSM mixing calculation

P. Boyle, N. Garron, J. Hudspith, A. Jüttner, **J. Kettle**, A. Khamseh, C. Lehner, A. Soni, JTT [1812.04981 PoS Lat'18, in preparation]

Flavour Physics and CKM: neutral meson mixing

$$\Delta m_P = |V_{tq_2}^* V_{tq_1}| \times f_P^2 m_P \hat{B}_P \times \text{known factors}$$



[HFLAV]

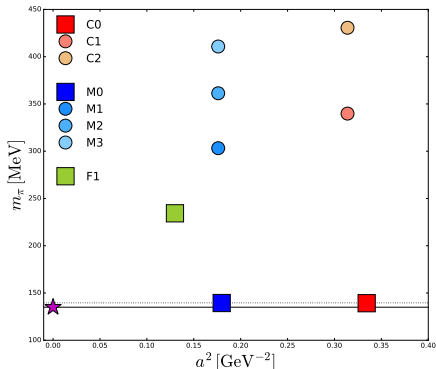
$$\Delta m_d = 0.5064 \pm 0.0019 \text{ ps}^{-1}$$

$$\Delta m_s = 17.757 \pm 0.021 \text{ ps}^{-1}$$

Computing ξ gives access to ratio V_{td}/V_{ts} :

$$\xi^2 = \frac{f_{B_s}^2 B_{B_s}}{f_B^2 B_B} = \left| \frac{V_{td}}{V_{ts}} \right|^2 \frac{\Delta m_s}{\Delta m_d} \frac{m_B}{m_{B_s}}$$

RBC/UKQCD $N_f = 2 + 1$ ensembles



- Iwasaki gauge action
- Domain Wall Fermion action
 - $\Rightarrow N_f = 2 + 1$ flavours in the sea
 - $\Rightarrow M_5 = 1.8$ for light and strange
- **2 ensembles with physical pion masses** [PRD 93 (2016) 074505]
- 3 Lattice spacings [JHEP 12 (2017) 008]
- Heavier m_π ensembles guide small chiral extrapolation of F1*

Chiral Fermions:

- $\Rightarrow O(a)$ improved
- \Rightarrow Multiplicative renormalisation

* F1 properties under investigation but expect only minor effects

Lattice set-up I

Light and strange

- Unitary light quark mass
- Physical strange quark mass
- DWF parameters same between sea and valence
- Gaussian source (sink) smearing for better overlap with ground state

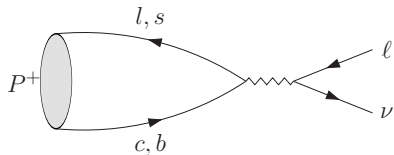
Heavy (charm and beyond)

- Möbius DWF
- $M_5 = 1.0$, $L_5 = 12$
- Stout smeared (3 hits, $\rho = 0.1$)
- Range of quark masses from below charm to $\sim m_b/2$ on finest ensemble

- ⇒ **All DWF** mixed action set-up
- ⇒ \mathbb{Z}_2 -noise sources (volume average) on every 2nd time slice
- ⇒ Increased heavy quark reach compared to [JHEP 04 (2016) 037, JHEP 12 (2017) 008]
- extrapolation towards b

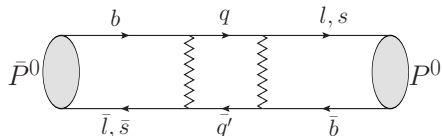
Lattice setup II

Leptonic decays:

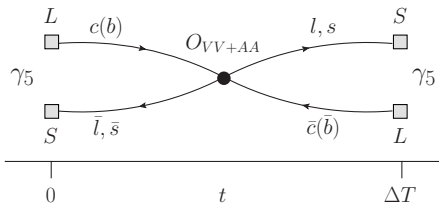
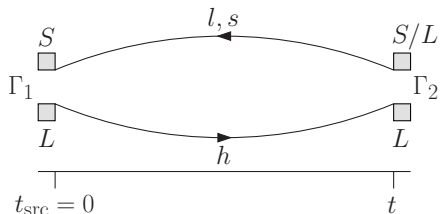


$$\mathcal{Z}_A \langle 0 | \bar{c} \gamma_4 \gamma_5 q | D_q(0) \rangle = f_{D_q} m_{D_q}$$

$P^0 - \bar{P}^0$ -mixing



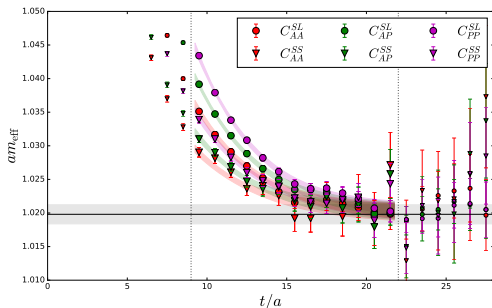
$$B_P = \frac{\langle \bar{P}^0 | O_{VV+AA} | P^0 \rangle}{8/3 f_P^2 m_P^2}$$



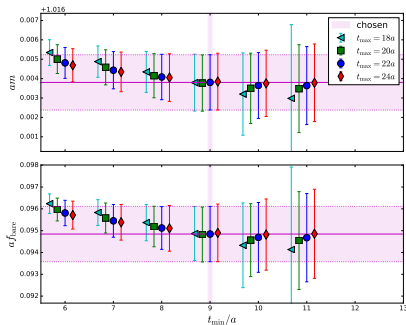
Many source-sink separations ΔT for 4-quark operator

Correlator Fitting - two point functions

Simultaneous two-exponential fit of 6 channels to extract masses and matrix elements of interest

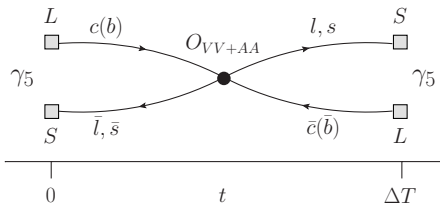
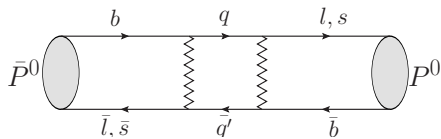


Example fit of worst case:
heavy-light meson with $am_h = 0.68$ on
M0



Stability

Correlator Fitting of 4-quark operators I



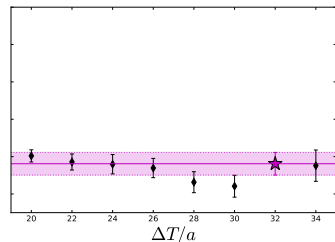
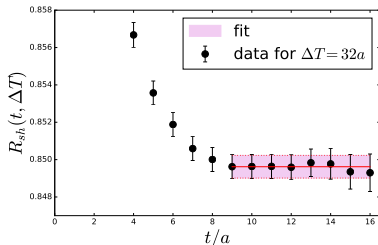
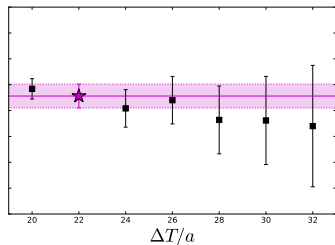
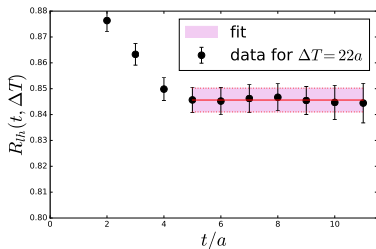
$$C_3(t, \Delta T) \equiv \langle P(\Delta T) O_{VV+AA}(t) \bar{P}^\dagger(0) \rangle$$

$$R(t, \Delta T) = \frac{C_3(t, \Delta T)}{8/3 C_{PA}(\Delta T - t) C_{AP}(t)} \rightarrow B_P \quad \text{for } t, \Delta T \gg 0$$

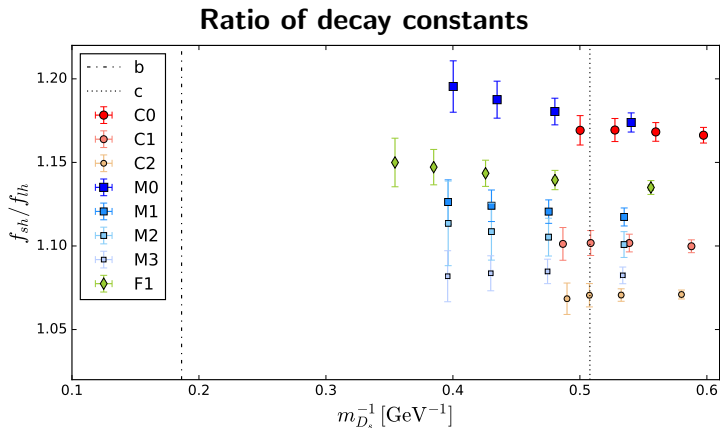
- Expect $R(t, \Delta T)$ to plateau for large t
- Check stability of plateau value by varying ΔT

Correlator Fitting of 4-quark operators II

Ex: $am_h = 0.68$ on M0

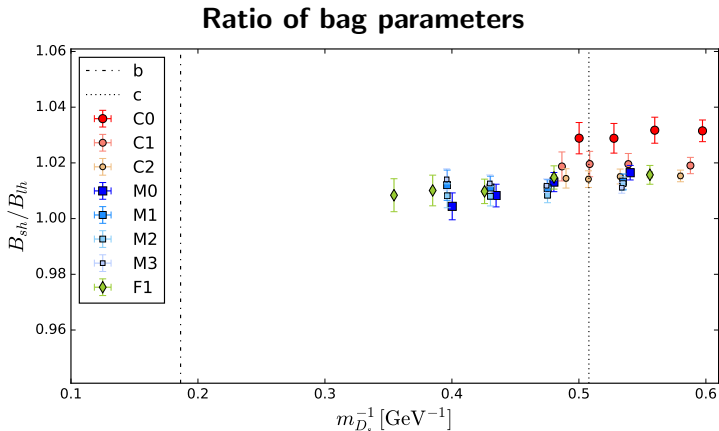


Results of correlator fits



- ⇒ Renormalisation constants cancel
- ⇒ Mild linear behaviour with $1/m_H$ and a^2
- ⇒ Stat precision: 0.4 - 1.0 %

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Global fit form

Base fit

$$O(a, m_\pi, m_H) = O(0, m_\pi^{\text{phys}}, m_H^{\text{phys}}) + C_{CL} a^2 + C_\chi \Delta m_\pi^2 + C_H \Delta m_H^{-1}$$

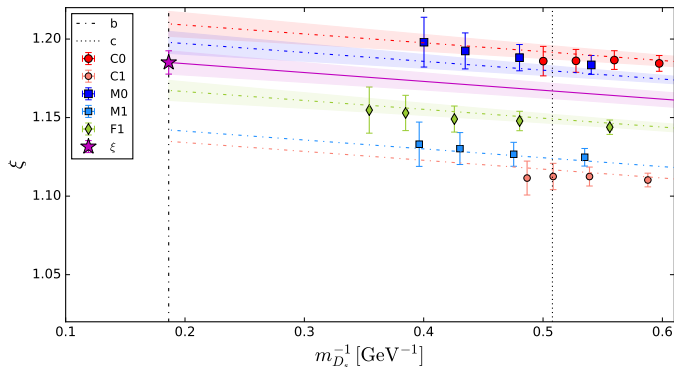
Assess systematic errors by

- varying cuts on pion mass
- using $m_H = m_D, m_{D_s}$ and m_{η_c}
- varying inclusion/exclusion of heaviest data points
- varying inclusion/exclusion of fit parameters
- including/estimating higher order terms ($a^4, (\Delta m_\pi^2)^2, (\Delta m_H^{-1})^2$)

⇒ Global fits are fully correlated.

Global fit results for ξ

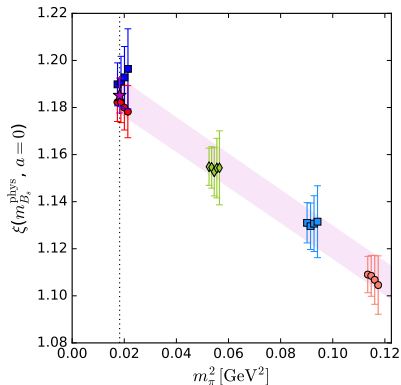
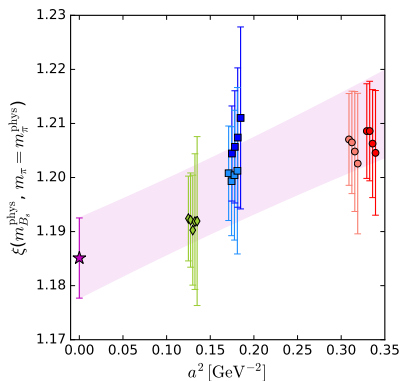
$$O(a, m_\pi, m_H) = O(0, m_\pi^{\text{phys}}, m_H^{\text{phys}}) + C_{CL} a^2 + C_\chi \Delta m_\pi^2 + C_H \Delta m_H^{-1}$$



Ratio of decay constants for $m_\pi \leq 350$ MeV

Global fit results for ξ

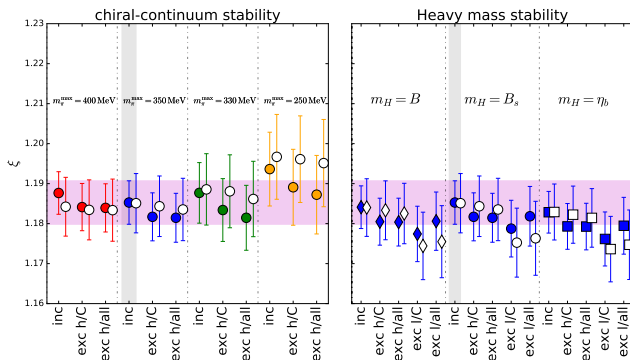
$$O(a, m_\pi, m_H) = O(0, m_\pi^{\text{phys}}, m_H^{\text{phys}}) + C_{CL}a^2 + C_\chi \Delta m_\pi^2 + C_H \Delta m_H^{-1}$$



Ratio of decay constants for $m_\pi \leq 350$ MeV

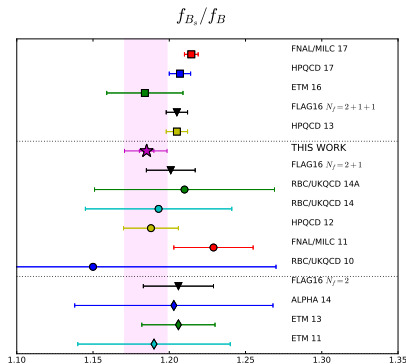
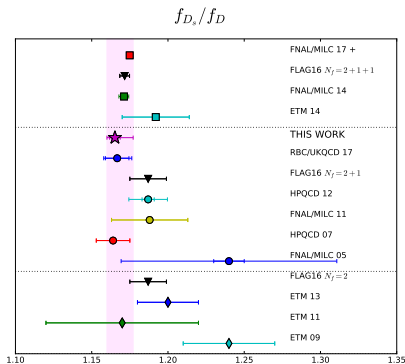
Systematic Errors - variations of cuts to data for ξ

- Global fits all correlated with satisfying p -values.
- sys error: includes chiral-CL (left), heavy mass (right), H.O. terms, $m_u \neq m_d$ and FV.



$$\xi = 1.1853(54)_{\text{stat}} \begin{pmatrix} +116 \\ -156 \end{pmatrix}_{\text{sys}}$$

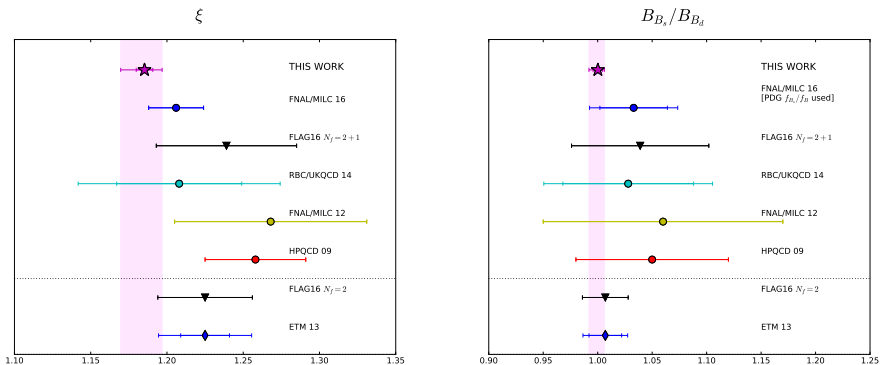
Comparison to literature - ratio of decay constants



- Self consistent with RBC/UKQCD17: JHEP **12** (2017) 008
- Complimentary to (most) literature - no effective action for b .
- One of few results with physical pion masses.

$$|V_{cd}/V_{cs}| = 0.2148(56)_{\text{exp}} \left(\begin{smallmatrix} +22 \\ -10 \end{smallmatrix} \right)_{\text{lat}}$$

Comparison to literature - ratio of mixing parameters



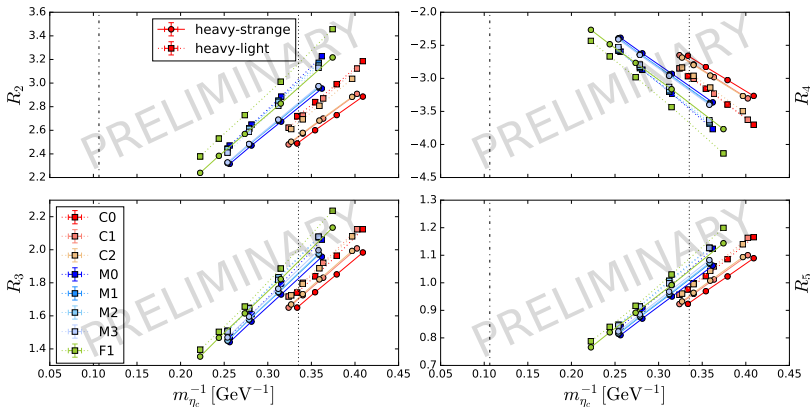
- Complimentary - no effective action needed for b
- Complimentary - **no operator mixing!**
- **First time with physical pion masses**

$$|V_{td}/V_{ts}| = 0.2018(4)_{\text{exp}} \left(\begin{matrix} +20 \\ -27 \end{matrix} \right)_{\text{lat}}$$

Next steps: Decay constants and bag parameters

- Different choice of (domain wall) action between light/strange and heavy quarks leads to a mixed action
- Mixed action renormalisation constants cancel for appropriate ratios (f_{B_s}/f_B , B_{B_s}/B_B), but are needed for individual decay constants and bag parameters.
- Need to carry out the fully non-perturbative mixed action renormalisation as outlined in JHEP **12** (2017) 008.
- Extend the study to the full BSM operator basis
⇒ analogous to RBC/UKQCD's $K - \bar{K}$ study (1812.04981, in preparation)

$B_{(s)}^0 - \bar{B}_{(s)}^0$ and $D^0 - \bar{D}^0$ PRELIMINARY and BARE



- “quite linear” in m_H^{-1}
- similar slopes for h-l and h-s
 $\Rightarrow SU(3)$ breaking rat's?

- renormalisation to be done
(mixed action + op mixing)
- analogous analysis to $K - \bar{K}$
paper + m_H dependence

Conclusions and Outlook

$SU(3)$ breaking ratios

- [arXiv:1812.08791](#)
- f_{D_s}/f_D , f_{B_s}/f_B , B_{B_s}/B_B and ξ
- $|V_{cd}/V_{cs}|$, $|V_{td}/V_{ts}|$
- 3 lattice spacings, 2 m_π^{phys}
- First result for ξ and B_{B_s}/B_B with m_π^{phys}
- m_h from below m_c to $\sim m_b/2$
⇒ extrapolation to b for ratios
⇒ fully relativistic
- Good continuum scaling and self-consistent
- Competitive precision

Ongoing

- Mixed action renormalisation of bilinears and four quark operators underway
- First results look promising
- ⇒ Determine $f_{B_{(s)}}$, $f_{D_{(s)}}$
- ⇒ Extend to full mixing operator basis for $B_{(s)}$ and compute short distance part of D .

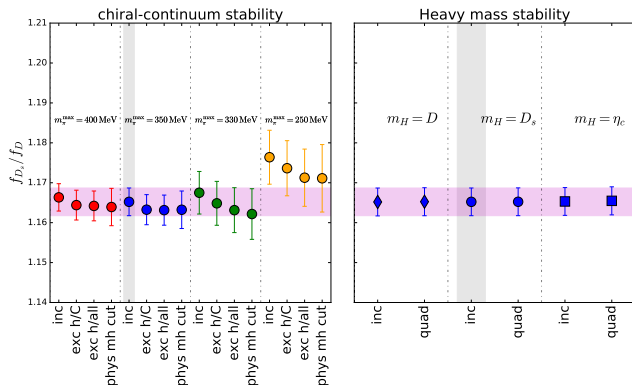
Outlook

- Supplement dataset with very fine JLQCD ensembles
- $a^{-1} = 2.8 \text{ GeV}$, $m_\pi = m_\pi^{\text{phys}}$

ADDITIONAL SLIDES

Systematic Errors - variations of cuts to data for f_{D_s}/f_D

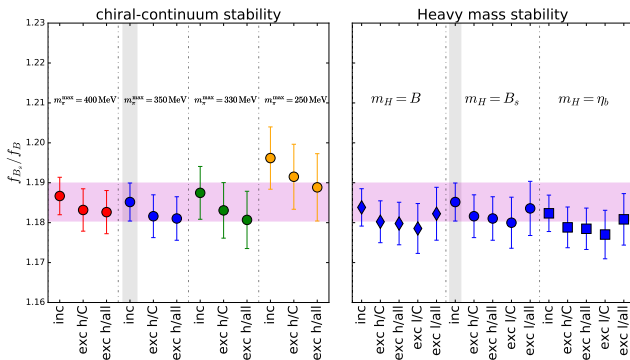
- Global fits all correlated with satisfying p -values.
- sys error: includes chiral-CL (left), heavy mass (right), H.O. terms, $m_u \neq m_d$ and FV.



$$f_{D_s}/f_D = 1.1652(35)_{\text{stat}} \begin{pmatrix} +120 \\ -52 \end{pmatrix}_{\text{sys}}$$

Systematic Errors - variations of cuts to data for f_{B_s}/f_B

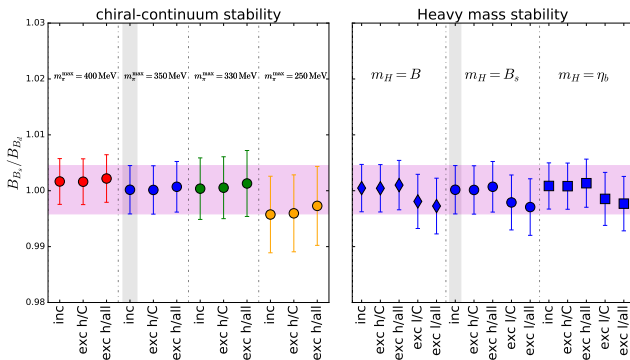
- Global fits all correlated with satisfying p -values.
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$$f_{B_s}/f_B = 1.1852(48)_{\text{stat}} \begin{pmatrix} +134 \\ -145 \end{pmatrix}_{\text{sys}}$$

Systematic Errors - variations of cuts to data for B_{B_s}/B_B

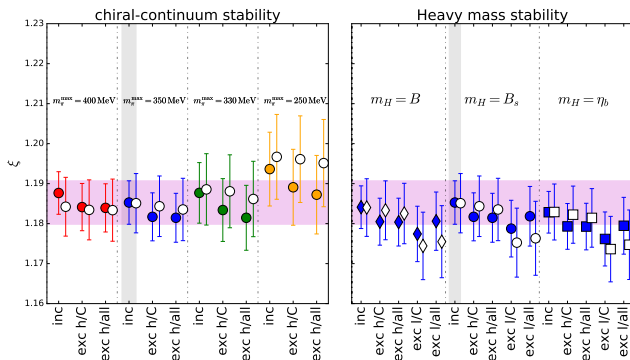
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$$B_{B_s}/B_B = 1.0002(43)_{\text{stat}} \left(\begin{matrix} +60 \\ -82 \end{matrix} \right)_{\text{sys}}$$

Systematic Errors - variations of cuts to data for ξ

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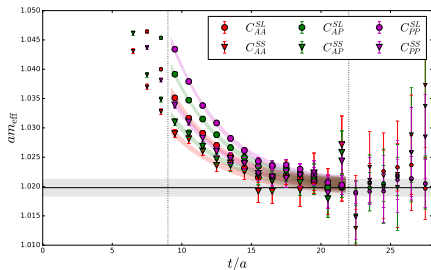
Cross checks of correlator fits I

$$C_{AP}^{LS}(t) \approx A_0^L P_0^S e^{-E_0 t} + A_1^L P_1^S e^{-E_1 t}$$

$$C_{AP}^{SS}(t) \approx A_0^S P_0^S e^{-E_0 t} + A_1^S P_1^S e^{-E_1 t}$$

Construct Linear Combination

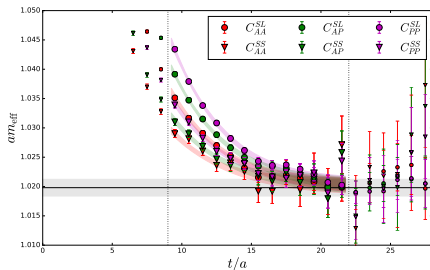
$$\begin{aligned} C_1^{AP}(t) &\equiv C_{AP}^{LS}(t) X^S - C_{AP}^{SS}(t) X^L \\ &\approx P_0^S \left(A_0^L X^S - A_0^S X^L \right) e^{-E_0 t} \\ &\quad + P_1^S \left(A_1^L X^S - A_1^S X^L \right) e^{-E_1 t} \end{aligned}$$



Cross checks of correlator fits I

$$C_1^{AP}(t) \approx P_0^S \left(A_0^L X^S - A_0^S X^L \right) e^{-E_0 t} \\ + P_1^S \underbrace{\left(A_1^L X^S - A_1^S X^L \right)}_{\text{small}} e^{-E_1 t}$$

Identify X^S, X^L with **central value** of A_1^S, A_1^L from fit.



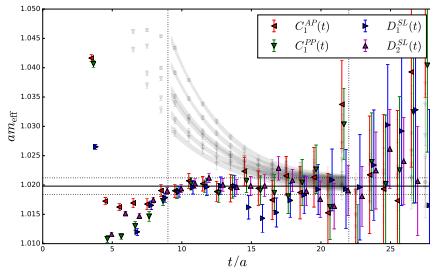
Cross checks of correlator fits I

$$C_1^{AP}(t) \approx P_0^S \left(A_0^L X^S - A_0^S X^L \right) e^{-E_0 t} \\ + P_1^S \underbrace{\left(A_1^L X^S - A_1^S X^L \right)}_{\text{small}} e^{-E_1 t}$$

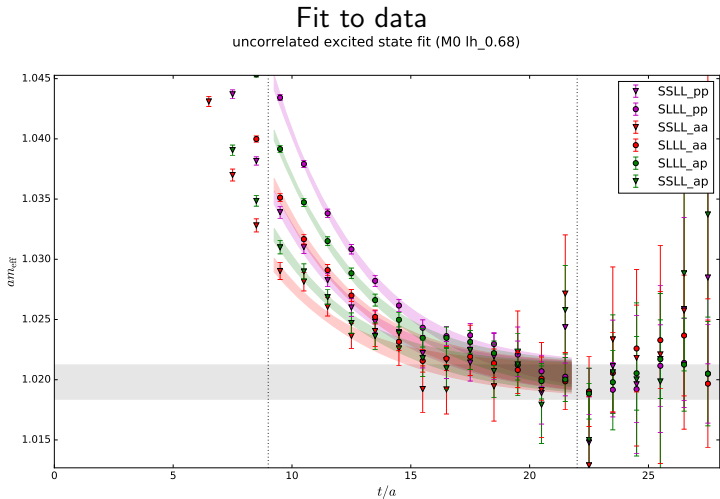
Identify X^S, X^L with **central value** of A_1^S, A_1^L from fit.

⇒ Removes (most of) excited state

⇒ Strong *a posteriori* check of fit range



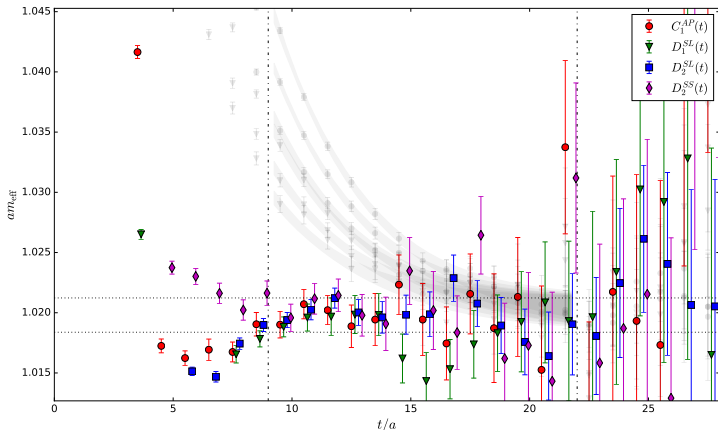
Cross checks of correlator fits II



Cross checks of correlator fits II

LCs using the central value of the fit result

Linear Combinations (M0 lh_0.68)



LCs plateau in fitrange region. \Rightarrow Excited state contamination removed.

Non-Perturbative Renormalisation of mixed action

SMOM ren. cond. relates amputated vertex functions to Z factors.

$$\begin{aligned} 1 &= \lim_{\bar{m} \rightarrow 0} \frac{1}{12q^2} \text{Tr} \left[(q \cdot \Lambda_A^{\text{ren}}) \gamma_5 \not{q} \right] \Big|_{\text{sym}} \\ &= \frac{Z_A}{Z_q} \lim_{\bar{m} \rightarrow 0} \frac{1}{12q^2} \text{Tr} \left[(q \cdot \Lambda_A^{\text{bare}}) \gamma_5 \not{q} \right] \Big|_{\text{sym}} \\ &\equiv \frac{Z_A}{Z_q} \mathcal{P}[\Lambda_A^{\text{bare}}] \end{aligned}$$

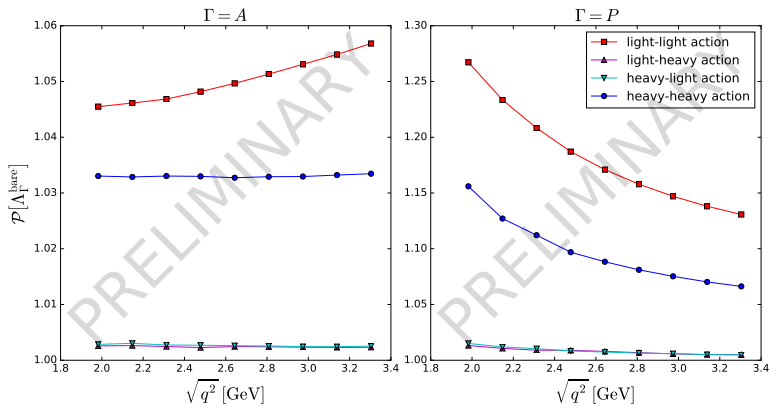
So for actions i, j

$$\frac{\mathcal{P}[\Lambda_A^{\text{bare}}]^{ii} \mathcal{P}[\Lambda_A^{\text{bare}}]^{jj}}{(\mathcal{P}[\Lambda_A^{\text{bare}}]^{ij})^2} = \frac{(Z_A^{ij})^2}{Z_A^{ii} Z_A^{jj}}$$

But for non-mixed actions we can determine Z_A^{ii} from conserved current.

Preliminary mixed action renormalisation

First study on single configuration



⇒ mixed NPR is feasible

⇒ need to compute Z_A^{hh} from conserved current to obtain Z_A^{hl}