

$\Delta b = 2$ mixings and ξ

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for the RBC-UKQCD Collaborations

Based on arXiv:1812.08791

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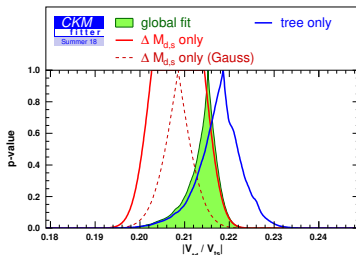
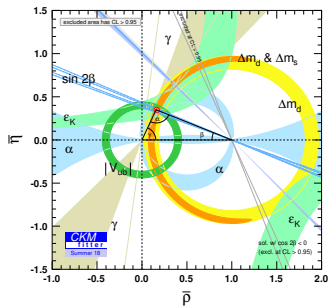


Outline

- 1 Introduction
- 2 Results for SU(3) breaking ratios (**arXiv:1812.08791**)
- 3 Ongoing Work
- 4 Conclusion and Outlook

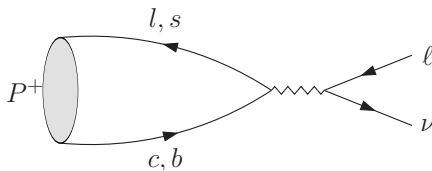
Motivation for charm and bottom flavour physics

- Huge experimental efforts: LHC, Belle II, BES III, ...
- Constrain CKM unitarity by combining non-perturbative input with experimental data.
- Test CKM matrix by determining the same CKM matrix element from different processes
- Constrain BSM models
- Address lepton flavour universality (violations?)



Flavour Physics and CKM: leptonic decay constants

Experiment \approx CKM \times Lattice \times (PT+kinematics)



Leptonic decays: $\Gamma(P \rightarrow l\nu_\ell) \approx |V_{q_2 q_1}|^2 \times f_P^2 \times \text{known factors}$

where $\mathcal{Z}_A \langle 0 | \bar{c} \gamma_4 \gamma_5 q | D_q(0) \rangle = f_{D_q} m_{D_q}$, $q = d, s$

[HFLAV+BESIII] $f_D |V_{cd}| = (45.9 \pm 1.1) \text{ MeV}$, $f_{D_s} |V_{cs}| = (249.1 \pm 3.2) \text{ MeV}$

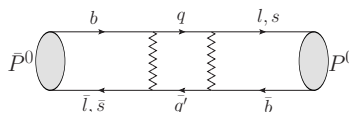
Computing f_{D_s}/f_D gives access to V_{cs}/V_{cd}

Neutral meson mixing

Neutral mesons oscillate with their antiparticles:

⇒ Difference between mass eigenstates: Δm^{exp} measured to $< 1\%$!

$$\Delta m \propto \overbrace{\langle B_{(s)}^0 | \mathcal{H}^{\Delta b=2} | \bar{B}_{(s)}^0 \rangle}^{\text{Short distance}} + \overbrace{\sum_n \frac{\langle B_{(s)}^0 | \mathcal{H}^{\Delta b=1} | n \rangle \langle n | \mathcal{H}^{\Delta b=1} | \bar{B}_{(s)}^0 \rangle}{E_n - M_{B_{(s)}}}}^{\text{Long distance}}$$



$$\propto \left| \sum_{q=u,c,t} \frac{m_q^2}{M_W^2} V_{qb} V_{ql}^* \right|^2 \approx \frac{m_t^4}{M_W^4} |V_{tb} V_{tl}^*|^2$$

SD: Top enhanced: $m_t^2 V_{tb} V_{tl}^* \gg m_c^2 V_{cb} V_{cl}^* \gg m_u^2 V_{ub} V_{ul}^*$

LD: Only m_c, m_u in intermediate states: no top + CKM suppressed

⇒ **Short distance dominated.**

Operator Product Expansion

Two scale problem: $\Lambda_{\text{QCD}} \sim 1 \text{ GeV} \ll m_{EW} \sim 100 \text{ GeV}$:

\Rightarrow Factorise via OPE

$$\Delta m \propto \sum_i C_i(\mu) \langle B_{(s)}^0 | \mathcal{O}_i^{\Delta b=2}(\mu) | \bar{B}_{(s)}^0 \rangle$$

- Perturbative model-dependent Wilson coefficients $C_i(\mu)$
- Non-perturbative model-independent matrix elements of $\mathcal{O}_i^{\Delta b=2}(\mu)$
- 5 independent (parity even) operators \mathcal{O}_i .

\Rightarrow SM: $\mathcal{O}_1 = (\bar{b}_a \gamma_\mu (\mathbb{1} - \gamma_5) q_a) (\bar{b}_b \gamma_\mu (\mathbb{1} - \gamma_5) q_b) = \mathcal{O}_{VV+AA}$
+ 4 (B)SM operators: $\mathcal{O}_2 - \mathcal{O}_5$

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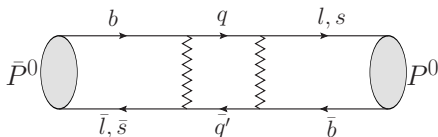
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RBC/UKQCD's $K - \bar{K}$ BSM mixing calculation

P. Boyle, N. Garron, J. Hudspith, A. Jüttner, **J. Kettle**, A. Khamseh, C. Lehner, A. Soni, JTT [1812.04981 PoS Lat'18, in preparation]

Flavour Physics and CKM: neutral meson mixing

$$\Delta m_P = |V_{tq_2}^* V_{tq_1}| \times f_P^2 m_P \hat{B}_P \times \text{known factors}$$



[HFLAV]

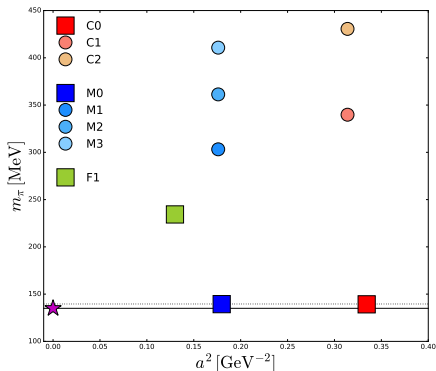
$$\Delta m_d = 0.5064 \pm 0.0019 \text{ ps}^{-1}$$

$$\Delta m_s = 17.757 \pm 0.021 \text{ ps}^{-1}$$

Computing ξ gives access to ratio V_{td}/V_{ts} :

$$\xi^2 \equiv \frac{f_{B_s}^2 B_{B_s}}{f_B^2 B_B} = \left| \frac{V_{td}}{V_{ts}} \right|^2 \frac{\Delta m_s}{\Delta m_d} \frac{m_B}{m_{B_s}}$$

RBC/UKQCD $N_f = 2 + 1$ ensembles



- Iwasaki gauge action
- Domain Wall Fermion action
 - $\Rightarrow N_f = 2 + 1$ flavours in the sea
 - $\Rightarrow M_5 = 1.8$ for light and strange
- **2 ensembles with physical pion masses** [PRD 93 (2016) 074505]
- 3 Lattice spacings [JHEP 12 (2017) 008]
- Heavier m_π ensembles guide small chiral extrapolation of F1

Chiral Fermions:

- $\Rightarrow O(a)$ improved
- \Rightarrow Multiplicative renormalisation

Lattice set-up

Light and strange

- Unitary light quark mass
- Physical strange quark mass
- DWF parameters same between sea and valence
- Gaussian source (sink) smearing for better overlap with ground state

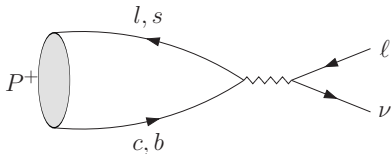
Heavy (charm and beyond)

- Möbius DWF
- $M_5 = 1.0$, $L_5 = 12$
- Stout smeared (3 hits, $\rho = 0.1$)
- Range of quark masses from below charm to $\sim m_b/2$ on finest ensemble

- ⇒ **All DWF** mixed action set-up
- ⇒ \mathbb{Z}_2 -noise sources (volume average) on every 2nd time slice
- ⇒ Increased heavy quark reach compared to [\[JHEP 04 \(2016\) 037, JHEP 12 \(2017\) 008\]](#)
- extrapolation towards b

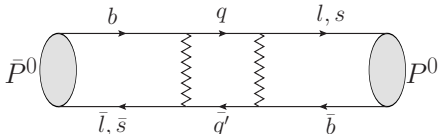
Measurement strategy

Leptonic decays:

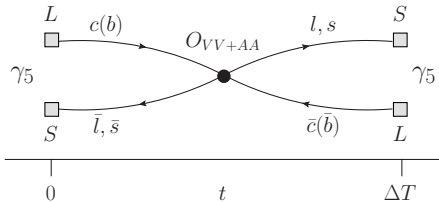
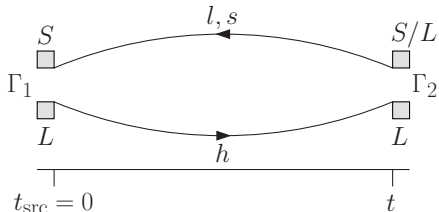


$$\mathcal{Z}_A \langle 0 | \bar{c} \gamma_4 \gamma_5 q | D_q(0) \rangle = f_{D_q} m_{D_q}$$

$P^0 - \bar{P}^0$ -mixing

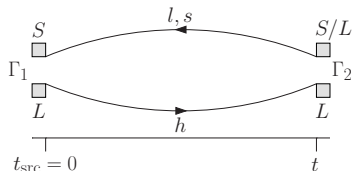


$$B_P = \frac{\langle \bar{P}^0 | O_{VV+AA} | P^0 \rangle}{8/3 f_P^2 m_P^2}$$



Many source-sink separations ΔT for 4-quark operator

Correlator fitting: strategy (2-point functions)



$$C_{ij}(t) = \sum_{n=0}^{\infty} (\psi_n)_i (\psi_n^*)_j e^{-E_n t}$$

with $E_n < E_{n+1}$ and $(\psi_n)_i = \frac{\langle 0 | O_i | n \rangle}{\sqrt{2E_n}}$ for $O = \bar{c}_2^L \Gamma q_1^X$ where $X = S, L$.

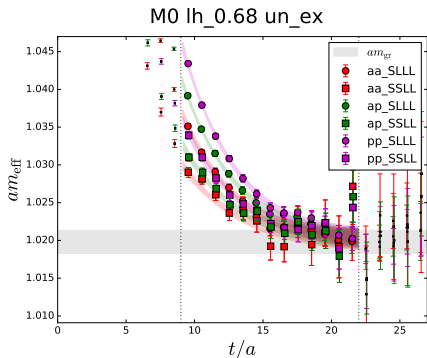
Consider $\Gamma = \gamma_5$ (**P**seudo scalar) and $\Gamma = \gamma_4 \gamma_5$ (**A**xial vector current).

ISSUE: Exponential noise growth i.e. **signal-to-noise problem**

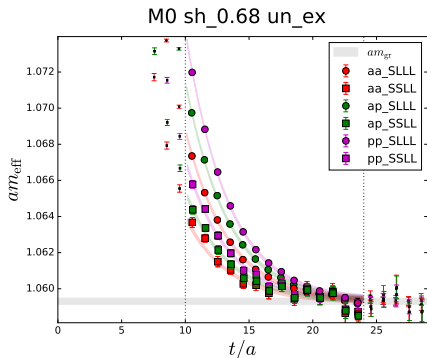
\Rightarrow Simultaneous uncorrelated excited state fits to 6 channels:

$\langle AA \rangle^{SL}$, $\langle AP \rangle^{SL}$, $\langle PP \rangle^{SL}$, $\langle AA \rangle^{SS}$, $\langle AP \rangle^{SS}$ and $\langle PP \rangle^{SS}$

Correlator fitting: fits (2-point functions)



Example fit (heavy-light meson with $am_h = 0.68$ on M0).



Example fit (heavy-strange meson with $am_h = 0.68$ on M0).

Correlator fitting: checks I

$$C_{AP}^{LS}(t) \approx A_0^L P_0^S e^{-E_0 t} + A_1^L P_1^S e^{-E_1 t}$$

$$C_{AP}^{SS}(t) \approx A_0^S P_0^S e^{-E_0 t} + A_1^S P_1^S e^{-E_1 t}$$

Construct Linear Combination

$$\begin{aligned} C_1^{AP}(t) &\equiv C_{AP}^{LS}(t) X^S - C_{AP}^{SS}(t) X^L \\ &\approx P_0^S \left(A_0^L X^S - A_0^S X^L \right) e^{-E_0 t} \\ &\quad + P_1^S \left(A_1^L X^S - A_1^S X^L \right) e^{-E_1 t} \end{aligned}$$

Correlator fitting: checks I

$$C_1^{AP}(t) \approx P_0^S \left(A_0^L X^S - A_0^S X^L \right) e^{-E_0 t} \\ + P_1^S \underbrace{\left(A_1^L X^S - A_1^S X^L \right)}_{\text{small}} e^{-E_1 t}$$

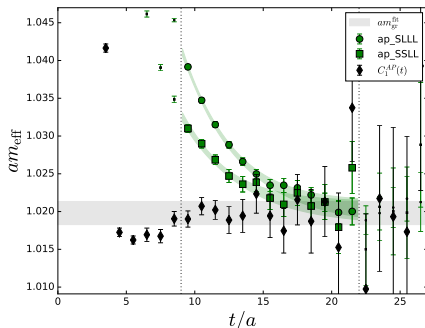
Identify X^S, X^L with **central value**
of A_1^S, A_1^L from fit.

Correlator fitting: checks I

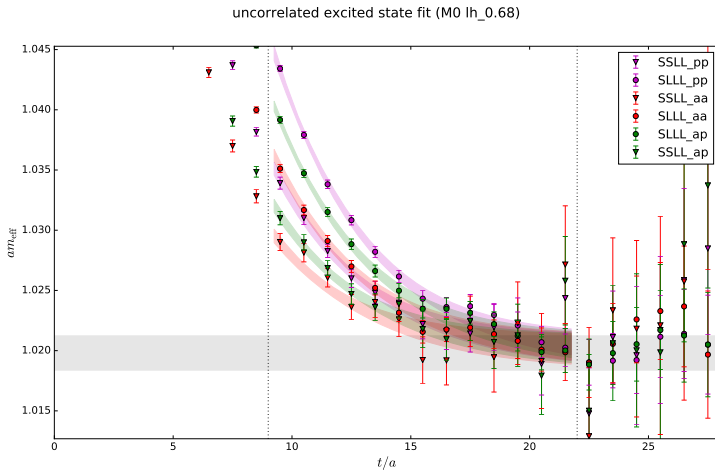
$$C_1^{AP}(t) \approx P_0^S \left(A_0^L X^S - A_0^S X^L \right) e^{-E_0 t} \\ + P_1^S \underbrace{\left(A_1^L X^S - A_1^S X^L \right)}_{\text{small}} e^{-E_1 t}$$

Identify X^S, X^L with **central value** of A_1^S, A_1^L from fit.

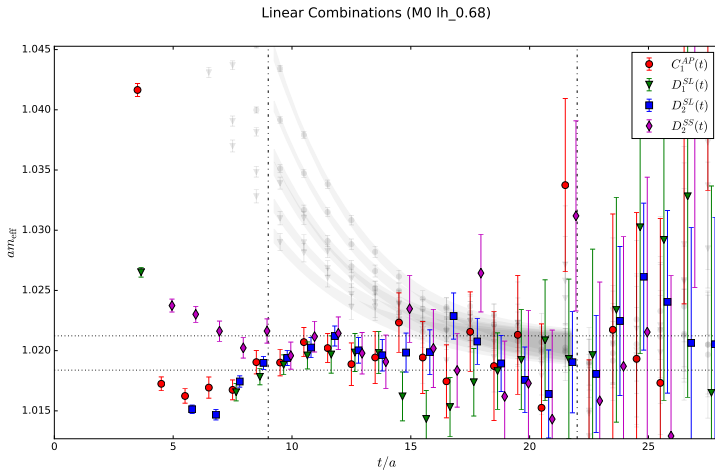
- ⇒ Removes (most of) excited state
- ⇒ Strong *a posteriori* check of fit range



Correlator fitting: checks II

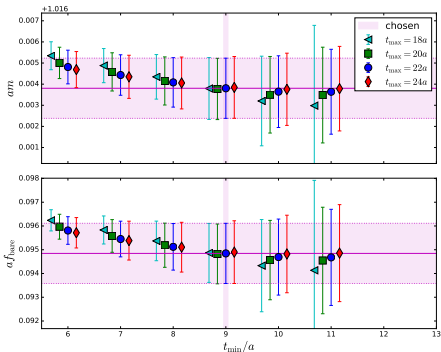


Correlator fitting: checks II

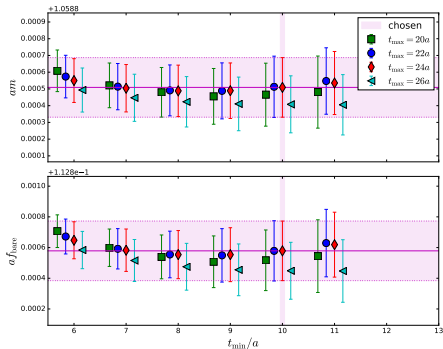


Correlator fitting: stability (2-point functions)

heavy-light on M0 ($am_h = 0.68$)

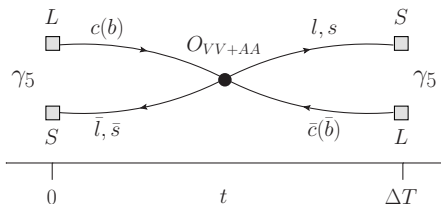
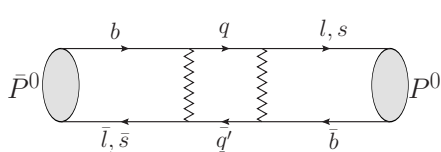


heavy-strange on M0 ($am_h = 0.68$)



Stability under variation of fit ranges

Correlator fitting of 4-quark operators: strategy



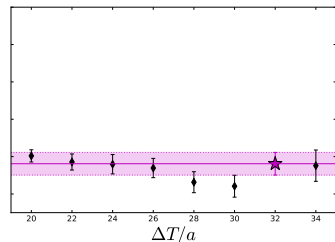
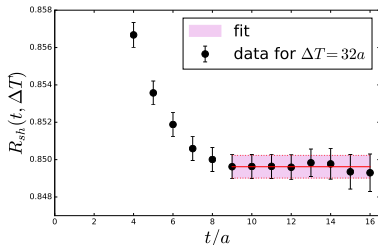
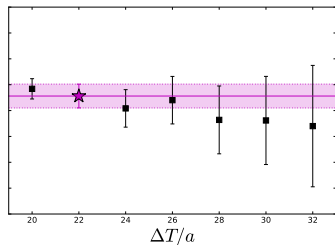
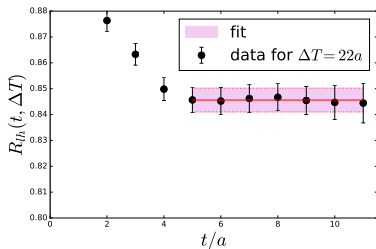
$$C_3(t, \Delta T) \equiv \langle P(\Delta T) O_{VV+AA}(t) \bar{P}^\dagger(0) \rangle$$

$$R(t, \Delta T) = \frac{C_3(t, \Delta T)}{8/3 C_{PA}(\Delta T - t) C_{AP}(t)} \rightarrow B_P \quad \text{for } t, \Delta T \gg 0$$

- Expect $R(t, \Delta T)$ to plateau for large t
- Check stability of plateau value by varying ΔT

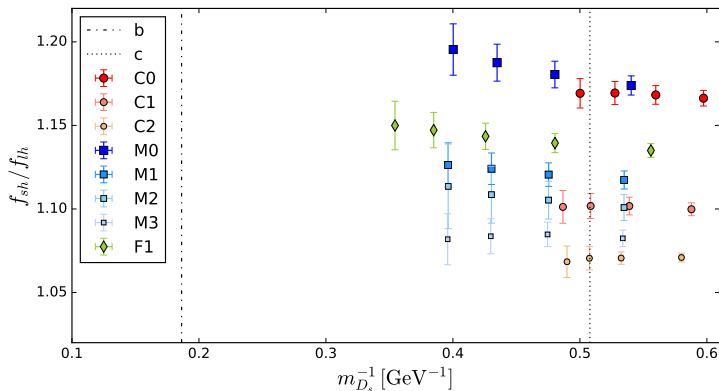
Correlator Fitting of 4-quark operators II

Ex: $am_h = 0.68$ on M0



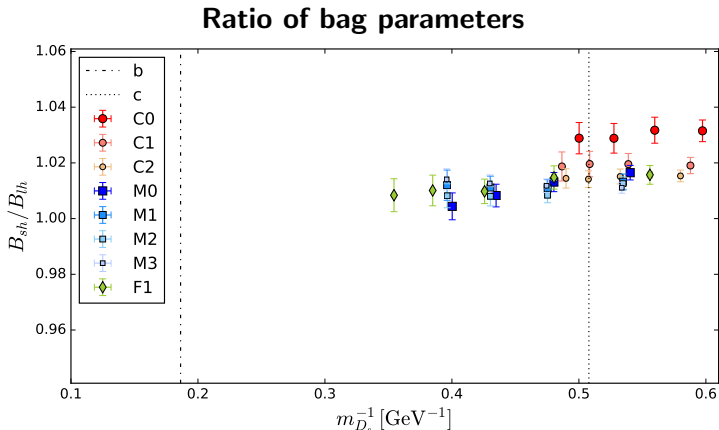
Results of correlator fits

Ratio of decay constants



- ⇒ Renormalisation constants cancel
- ⇒ Mild linear behaviour with $1/m_H$ and a^2
- ⇒ Stat precision: 0.4 - 1.0 %

Results of correlator fits



- ⇒ Renormalisation constants cancel
- ⇒ Mild linear behaviour with $1/m_H$ and a^2
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Global fit: ansatz

Base fit

$$O(a, m_\pi, m_H) = O(0, m_\pi^{\text{phys}}, m_H^{\text{phys}}) + C_{CL} a^2 + C_\chi \Delta m_\pi^2 + C_H \Delta m_H^{-1}$$

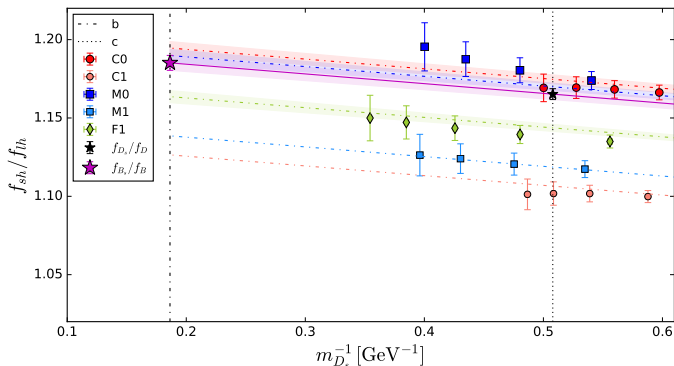
Assess systematic errors by

- varying cuts on pion mass
- using $m_H = m_D, m_{D_s}$ and m_{η_c}
- varying inclusion/exclusion of heaviest data points
- varying inclusion/exclusion of fit parameters
- including/estimating higher order terms ($a^4, (\Delta m_\pi^2)^2, (\Delta m_H^{-1})^2$)

⇒ Global fits are fully correlated.

Global fit: ratio of decay constants

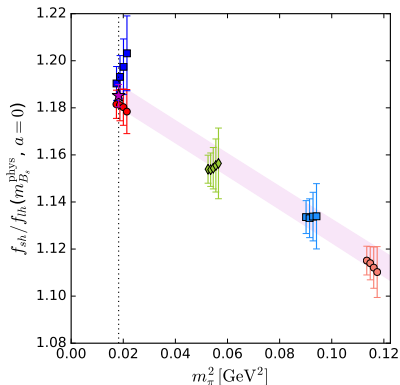
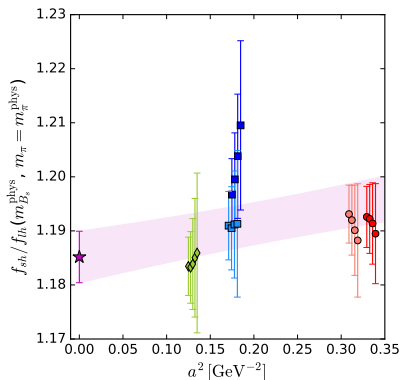
$$O(a, m_\pi, m_H) = \frac{f_{B_s}}{f_B} + C_{CL} a^2 + C_\chi \Delta m_\pi^2 + C_H \Delta m_H^{-1}$$



Ratio of decay constants for $m_\pi \leq 350$ MeV

Global fit: ratio of decay constants

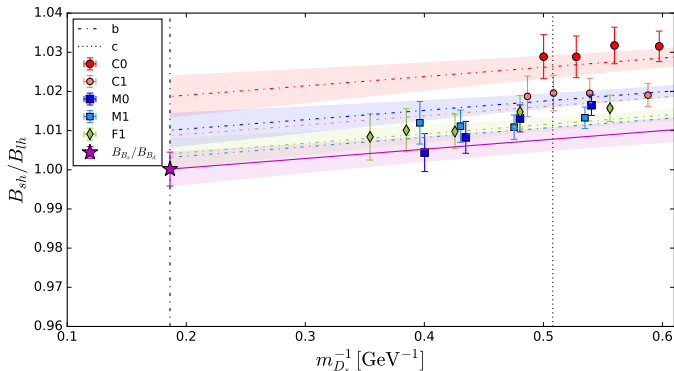
$$O(a, m_\pi, m_H) = \frac{f_{B_s}}{f_B} + C_{CL}a^2 + C_\chi \Delta m_\pi^2 + C_H \Delta m_H^{-1}$$



Ratio of decay constants for $m_\pi \leq 350$ MeV

Global fit: ratio of bag parameters

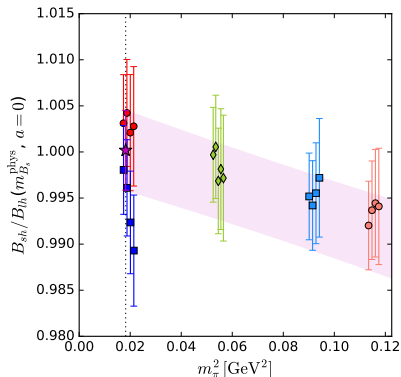
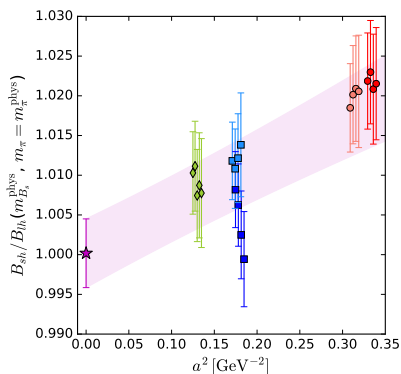
$$O(a, m_\pi, m_H) = \frac{B_{B_s}}{B_B} + C_{CL} a^2 + C_\chi \Delta m_\pi^2 + C_H \Delta m_H^{-1}$$



Ratio of bag parameters for $m_\pi \leq 350$ MeV

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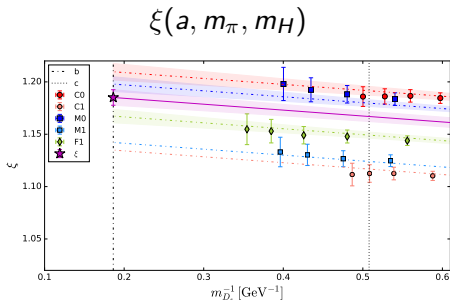
Ratio of bag parameters for $m_\pi \leq 350$ MeV

Global fit results - ratio of bag parameters and ξ

Recall:

$$\xi \equiv f_{B_s}/f_B \times \sqrt{B_{B_s}/B_B}$$

- 1 chiral-CL of product of ratios
- 2 product of chiral-CL of ratios.



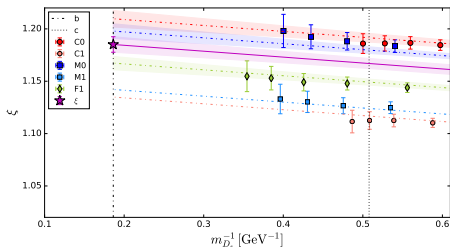
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$\xi(a, m_\pi, m_H)$



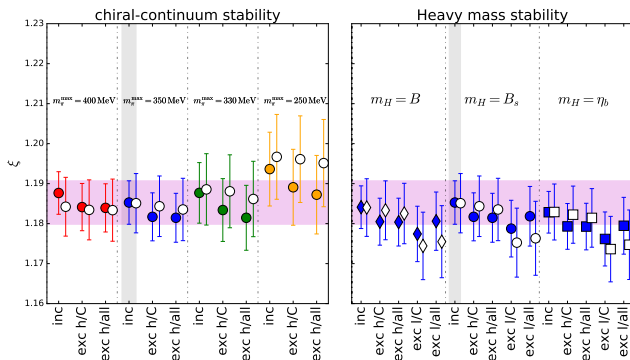
$$\lim_{a \rightarrow 0; m_q \rightarrow \text{phys}} \left[f_{hs}/f_{hl} \sqrt{B_{hs}/B_{hl}} \right] (a, m_\pi, m_H) = 1.1851(74)_{\text{stat}}$$

$$[f_{B_s}/f_B]_{\text{phys}} \times \sqrt{[B_{B_s}/B_B]_{\text{phys}}} = 1.1853(54)_{\text{stat}}$$

chiral continuum limit of individual ratios gives better signal

Systematic Errors - variations of cuts to data for ξ

- Global fits all correlated with satisfying p -values.
- sys error: includes chiral-CL (left), heavy mass (right), H.O. terms, $m_u \neq m_d$ and FV.



$$\xi = 1.1853(54)_{\text{stat}} \begin{pmatrix} +116 \\ -156 \end{pmatrix}_{\text{sys}}$$

Limitations and “ultimate precision”

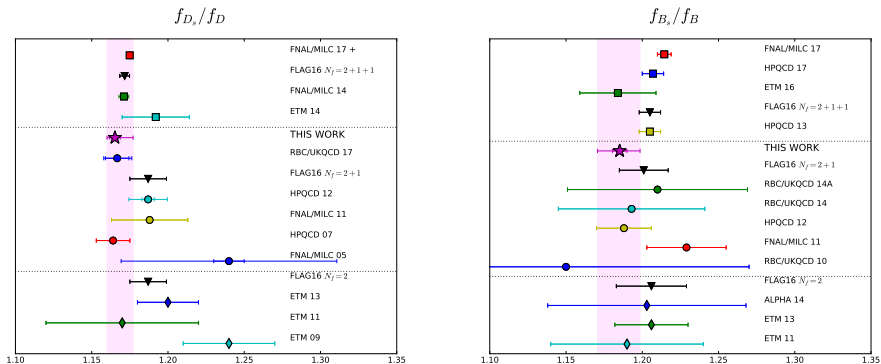
Experimental precision on $\Delta m_s \sim 0.1\%$ and $\Delta m_d \sim 0.4\%$.

Theoretical precision on $\xi \sim 1.3\%$

	f_{D_s}/f_D		f_{B_s}/f_B		ξ		B_{B_s}/B_{B_d}	
	absolute	relative	absolute	relative	absolute	relative	absolute	relative
central	1.1652		1.1852		1.1853		1.0002	
stat	0.0035	0.30%	0.0048	0.40%	0.0054	0.46%	0.0043	0.43%
fit chiral-CL	+0.0112 -0.0031	+0.96 % -0.26 %	+0.0110 -0.0045	+0.93 % -0.38 %	+0.0084 -0.0038	+0.71 % -0.32 %	+0.0020 -0.0044	+0.20 % -0.44 %
fit heavy mass	+0.0003 -0.0000	+0.02 % -0.00 %	+0.0000 -0.0081	+0.00 % -0.69 %	+0.0000 -0.0091	+0.00 % -0.77 %	+0.0012 -0.0031	+0.12 % -0.31 %
H.O. heavy	0.0000	0.00%	0.0054	0.45%	0.0049	0.41%	0.0021	0.21%
H.O. disc.	0.0009	0.07%	0.0009	0.07%	0.0021	0.18%	0.0016	0.16%
$m_u \neq m_d$	0.0009	0.08%	0.0009	0.07%	0.0010	0.08%	0.0001	0.01%
finite size	0.0021	0.18%	0.0021	0.18%	0.0021	0.18%	0.0018	0.18%
total systematic	+0.0114 -0.0039	+0.98 % -0.34 %	+0.0125 -0.0137	+1.06 % -1.16 %	+0.0102 -0.0146	+0.86 % -1.24 %	+0.0041 -0.0070	+0.41 % -0.70 %
total sys+stat	+0.0120 -0.0052	+1.03 % -0.45 %	+0.0134 -0.0145	+1.13 % -1.22 %	+0.0116 -0.0156	+0.97 % -1.32 %	+0.0060 -0.0082	+0.60 % -0.82 %

⇒ Systematically Improvable with finer lattices at (near) physical m_π .

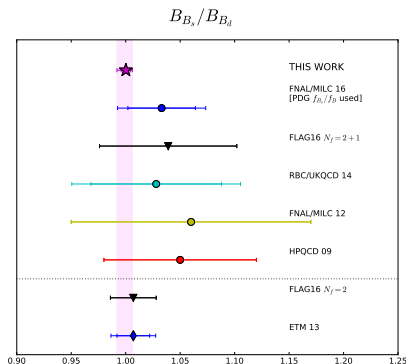
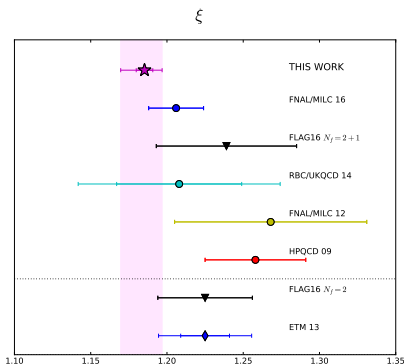
Comparison to literature - ratio of decay constants



- Self consistent with RBC/UKQCD17: JHEP **12** (2017) 008
- Complimentary to (most) literature - no effective action for b .
- One of few results with physical pion masses.

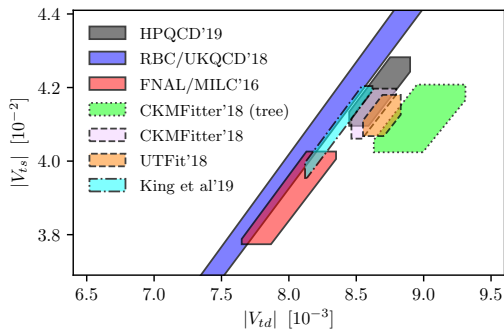
$$|V_{cd}/V_{cs}| = 0.2148(56)_{\text{exp}} \left(\begin{smallmatrix} +22 \\ -10 \end{smallmatrix} \right)_{\text{lat}}$$

Comparison to literature - ratio of mixing parameters



- Complimentary - no effective action needed for b
- Complimentary - **no operator mixing!**
- **First time with physical pion masses**
- New results: HPQCD'19, King et al. '19

Comparison to literature - V_{td}/V_{ts}



Plot taken from HPQCD'19

$$|V_{td}/V_{ts}| = 0.2018(4)_e \begin{pmatrix} +20 \\ -27 \end{pmatrix}_t$$

- Slight “discrepancy” between tree-only and loop determinations
- Error still dominated by theory
- Requires more work, but groups are active
- Our next target: V_{td} and V_{ts}

Next steps: Decay constants and bag parameters

- ① Different choice of (domain wall) action between light/strange and heavy quarks leads to a mixed action
Mixed action renormalisation constants cancel for appropriate ratios (f_{B_s}/f_B , B_{B_s}/B_B), but are needed for individual decay constants and bag parameters
- ⇒ Need to carry out the fully non-perturbative mixed action renormalisation as outlined in JHEP **12** (2017) 008.

Next steps: Decay constants and bag parameters

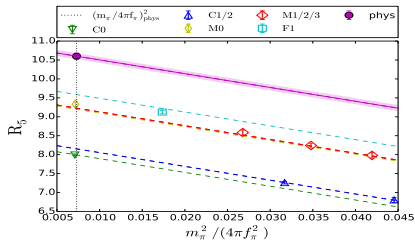
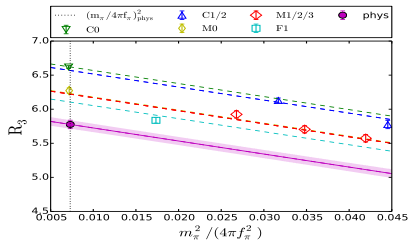
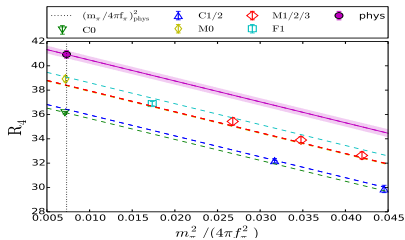
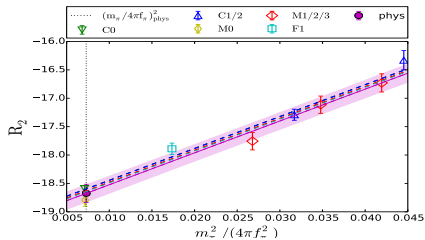
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- 3 Supplement data set with JLQCD ensemble (in collaboration with S. Hashimoto and T. Kaneko)
⇒ further reach in m_H due to finer lattice spacing

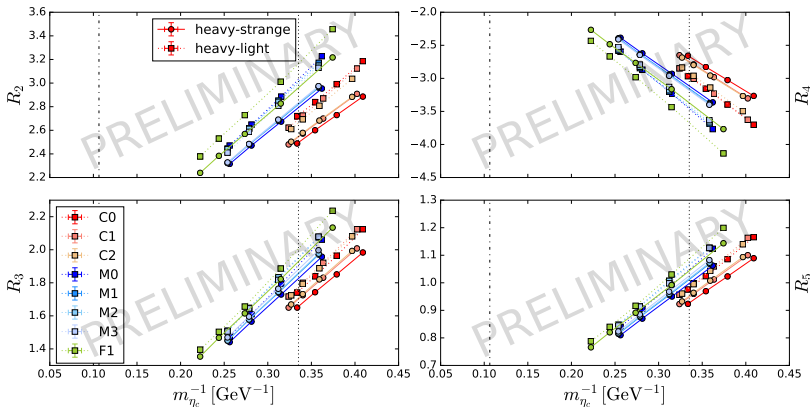
preliminary $K^0 - \bar{K}^0$ results [1812.04981, in preparation]

$$R_i \equiv \langle \bar{P}^0 | \mathcal{O}_i | P^0 \rangle / \langle \bar{P}^0 | \mathcal{O}_1 | P^0 \rangle$$



PRELIMINARY RESULTS in \overline{MS} at 3 GeV

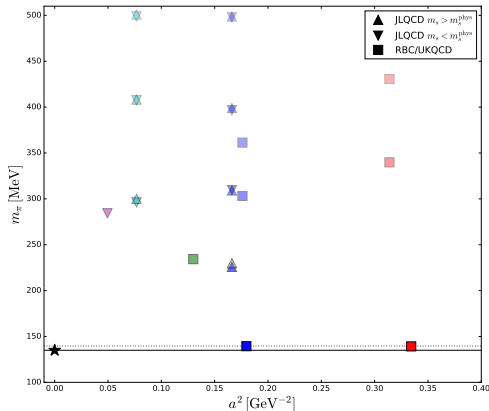
$B_{(s)}^0 - \bar{B}_{(s)}^0$ (and $D^0 - \bar{D}^0$) PRELIMINARY and BARE



- “quite linear” in m_H^{-1}
- similar slopes for h-l and h-s
 $\Rightarrow SU(3)$ breaking rat's?

- renormalisation to be done
 (mixed action + op mixing)
- analogous analysis to $K - \bar{K}$
 paper + m_H dependence

Increased set of ensembles



JLQCD (triangles)

Fine lattices:

$$a^{-1} = 2.4 - 4.5 \text{ GeV}$$

UKQCD (squares)

+RBC Physical Pion masses

Both: $N_f = 2 + 1$ DWF

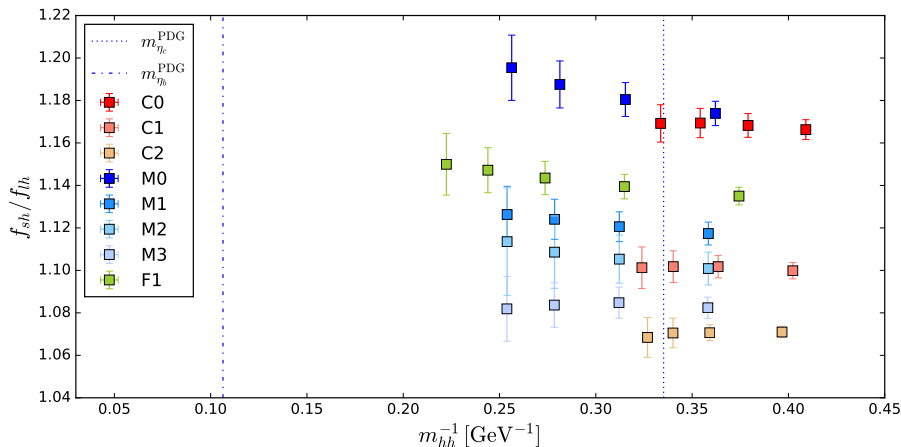
3+3 Lattice Spacings

⇒ Fine lattices: Further heavy quark reach on JLQCD ensembles

⇒ Chiral extrapolation stabilised by m_π^{phys} ensembles

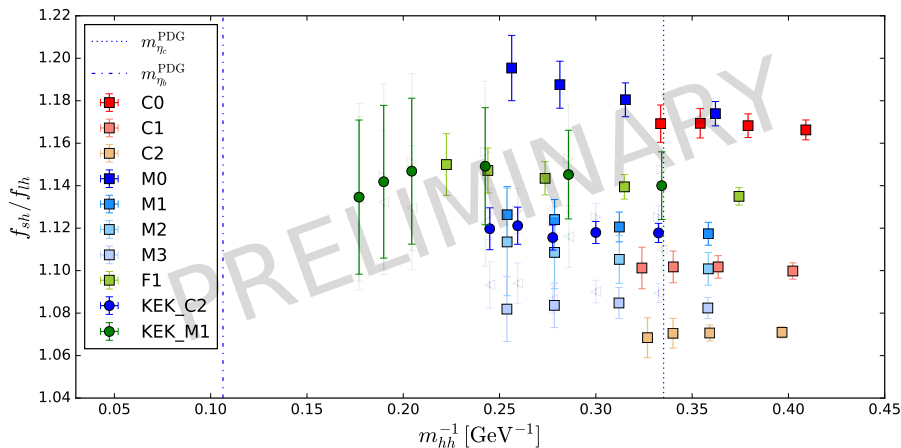
⇒ **Combined physics analysis with S. Hashimoto and T. Kaneko**

JLQCD + RBC/UKQCD data: ratio of decay constants I



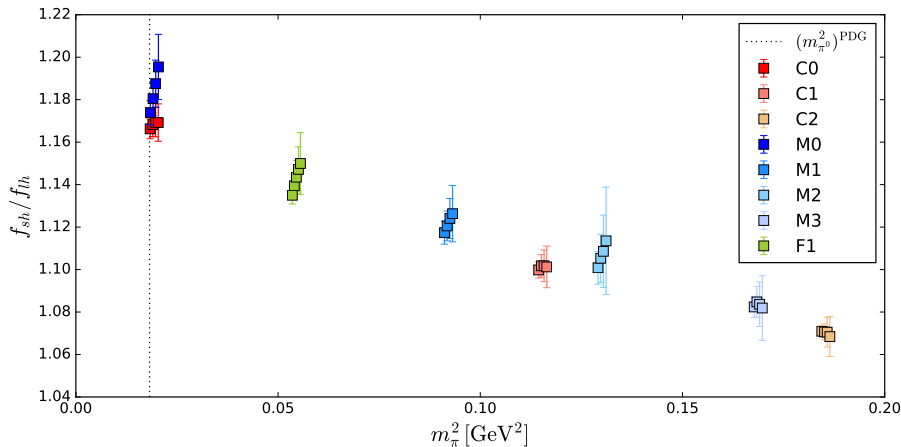
RBC-UKQCD data set from arXiv:1812.08791

JLQCD + RBC/UKQCD data: ratio of decay constants I

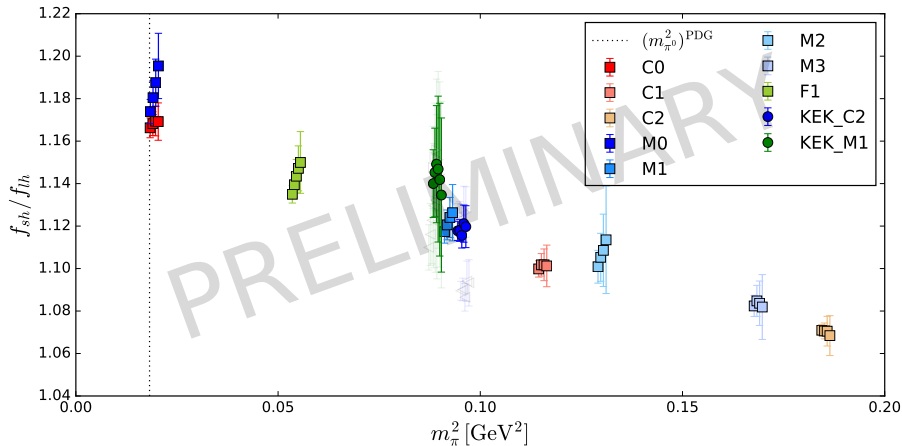


Increased reach in the heavy-mass. The fine KEK ensemble is yet to come!

JLQCD + RBC/UKQCD data: ratio of decay constants II



JLQCD + RBC/UKQCD data: ratio of decay constants II



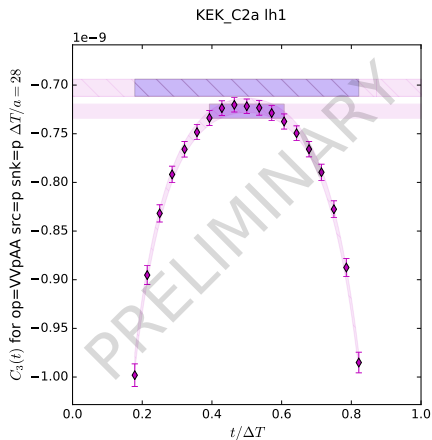
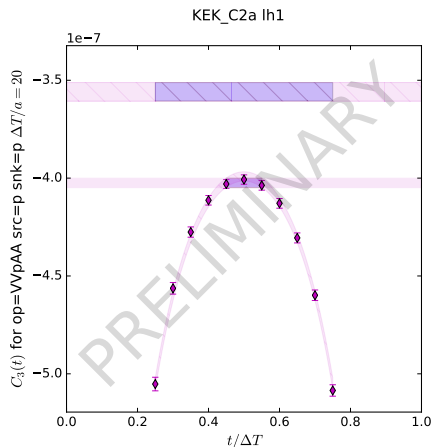
JLQCD + RBC/UKQCD data: bag parameter fit strategy

Write out the excited state contribution to the three point functions

$$\begin{aligned} C_3^{\mathcal{O}}(t; \Delta T) \approx & + \frac{M_{\text{snk}}^0 M_{\text{src}}^0}{4E_0 E_0} \langle gr | \mathcal{O} | gr \rangle e^{-E_0 \Delta T} \\ & + \frac{M_{\text{snk}}^0 M_{\text{src}}^1}{4E_0 E_1} \langle gr | \mathcal{O} | ex \rangle e^{-(E_0 + E_1) \Delta T / 2} e^{-(E_1 - E_0)(t - \Delta T / 2)} \\ & + \frac{M_{\text{snk}}^1 M_{\text{src}}^0}{4E_0 E_1} \langle ex | \mathcal{O} | gr \rangle e^{-(E_0 + E_1) \Delta T / 2} e^{-(E_0 - E_1)(t - \Delta T / 2)} \end{aligned}$$

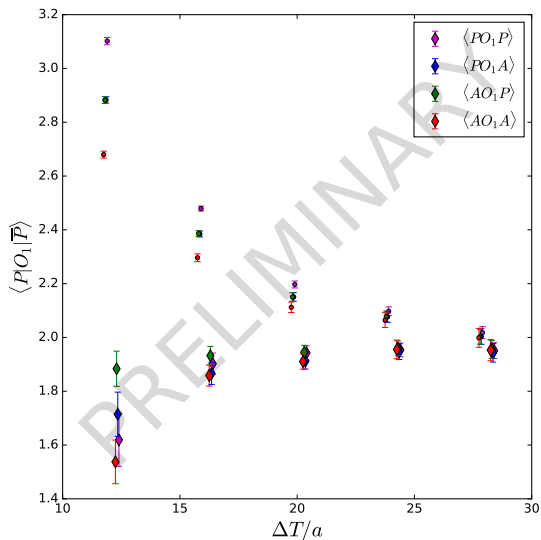
- Use the excited state information from the two-point fits.
- Compare ground state only vs. ground + excited as a function of ΔT

JLQCD + RBC/UKQCD data: bag parameter fitting



- Ground state fit: to a constant (range indicated in plot)
- Excited state fit: to functional form above for all data points shown.

JLQCD + RBC/UKQCD data: bag parameter fitting



- Combined fit to 2-point and 3-point functions possible
- Small symbols: ground state only
- Large symbols: excited state fit
- Each data point is a separate fit

- Stable for $\Delta T/a \gtrsim 20$
- Stable for different channels

Conclusions and Outlook

$SU(3)$ breaking ratios

- [arXiv:1812.08791](#)
- f_{D_s}/f_D , f_{B_s}/f_B , B_{B_s}/B_B and ξ
- $|V_{cd}/V_{cs}|$, $|V_{td}/V_{ts}|$
- 3 lattice spacings, 2 m_π^{phys}
- First result for ξ and B_{B_s}/B_B with m_π^{phys}
- m_h from below m_c to $\sim m_b/2$
 \Rightarrow extrapolation to b for ratios
 \Rightarrow fully relativistic
- Good continuum scaling and self-consistent
- Competitive precision

Ongoing

- Supplementing dataset with very fine JLQCD ensembles
- Combined fit with universality constraint
- Mixed action renormalisation underway
- First results look promising
- \Rightarrow Determine $f_{B_{(s)}}$, $f_{D_{(s)}}$
- \Rightarrow Full mixing operator basis for $B_{(s)}$ and D (short distance).

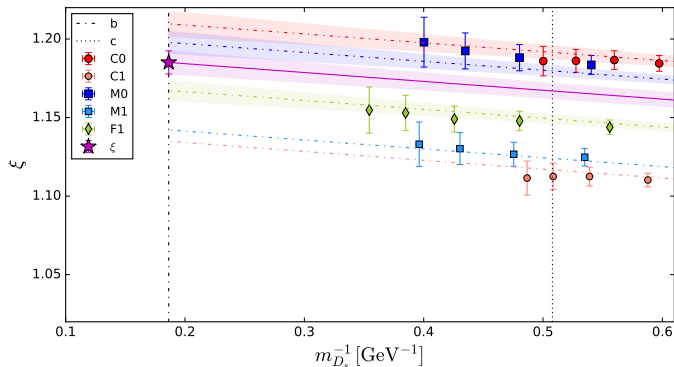
Outlook

- $a^{-1} = 2.8 \text{ GeV}$, $m_\pi = m_\pi^{\text{phys}}$

ADDITIONAL SLIDES

Global fit results for ξ

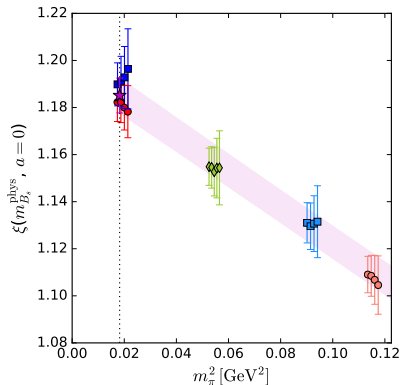
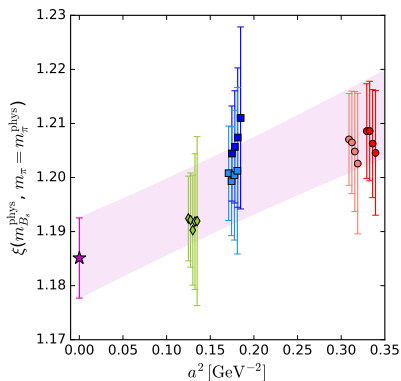
$$O(a, m_\pi, m_H) = O(0, m_\pi^{\text{phys}}, m_H^{\text{phys}}) + C_{CL} a^2 + C_\chi \Delta m_\pi^2 + C_H \Delta m_H^{-1}$$



Ratio of decay constants for $m_\pi \leq 350$ MeV

Global fit results for ξ

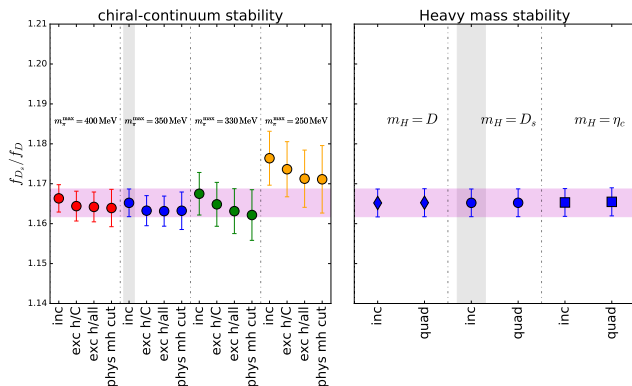
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Ratio of decay constants for $m_\pi \leq 350 \text{ MeV}$

Systematic Errors - variations of cuts to data for f_{D_s}/f_D

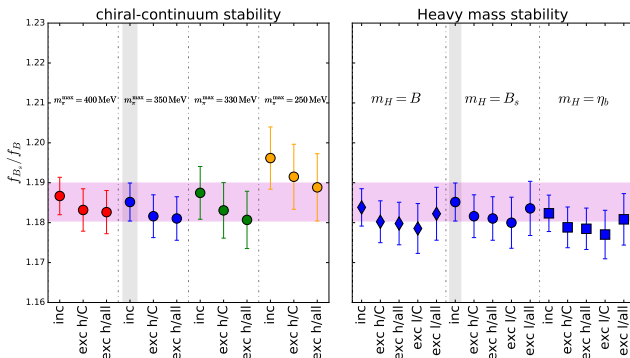
- Global fits all correlated with satisfying p -values.
- sys error: includes chiral-CL (left), heavy mass (right), H.O. terms, $m_u \neq m_d$ and FV.



$$f_{D_s}/f_D = 1.1652(35)_{\text{stat}} \left(\begin{matrix} +120 \\ -52 \end{matrix} \right)_{\text{sys}}$$

Systematic Errors - variations of cuts to data for f_{B_s}/f_B

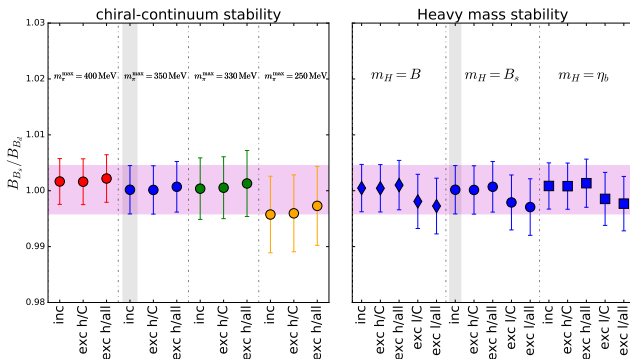
- Global fits all correlated with satisfying p -values.
- sys error: includes chiral-CL (left), heavy mass (right), H.O. terms, $m_u \neq m_d$ and FV.



$$f_{B_s}/f_B = 1.1852(48)_{\text{stat}} \left(\begin{matrix} +134 \\ -145 \end{matrix} \right)_{\text{sys}}$$

Systematic Errors - variations of cuts to data for B_{B_s}/B_B

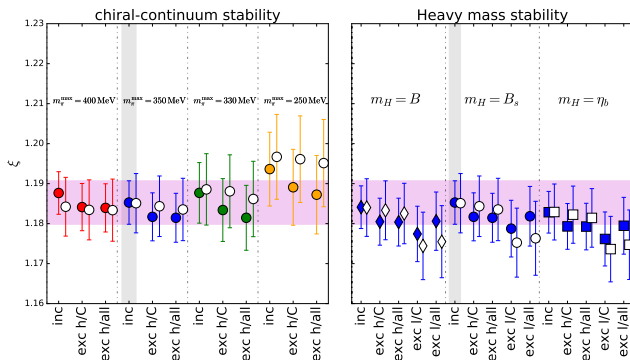
- Global fits all correlated with satisfying p -values.
- sys error: includes chiral-CL (left), heavy mass (right), H.O. terms, $m_u \neq m_d$ and FV.



$$B_{B_s}/B_B = 1.0002(43)_{\text{stat}} \left(\begin{matrix} +60 \\ -82 \end{matrix} \right)_{\text{sys}}$$

Systematic Errors - variations of cuts to data for ξ

- Global fits all correlated with satisfying p -values.
- sys error: includes chiral-CL (left), heavy mass (right), H.O. terms, $m_u \neq m_d$ and FV.



$$\xi = 1.1853(54)_{\text{stat}} \begin{pmatrix} +116 \\ -156 \end{pmatrix}_{\text{sys}}$$

Non-Perturbative Renormalisation of mixed action

SMOM ren. cond. relates amputated vertex functions to Z factors.

$$\begin{aligned} 1 &= \lim_{\bar{m} \rightarrow 0} \frac{1}{12q^2} \text{Tr} \left[(q \cdot \Lambda_A^{\text{ren}}) \gamma_5 \not{q} \right] |_{\text{sym}} \\ &= \frac{Z_A}{Z_q} \lim_{\bar{m} \rightarrow 0} \frac{1}{12q^2} \text{Tr} \left[(q \cdot \Lambda_A^{\text{bare}}) \gamma_5 \not{q} \right] |_{\text{sym}} \\ &\equiv \frac{Z_A}{Z_q} \mathcal{P}[\Lambda_A^{\text{bare}}] \end{aligned}$$

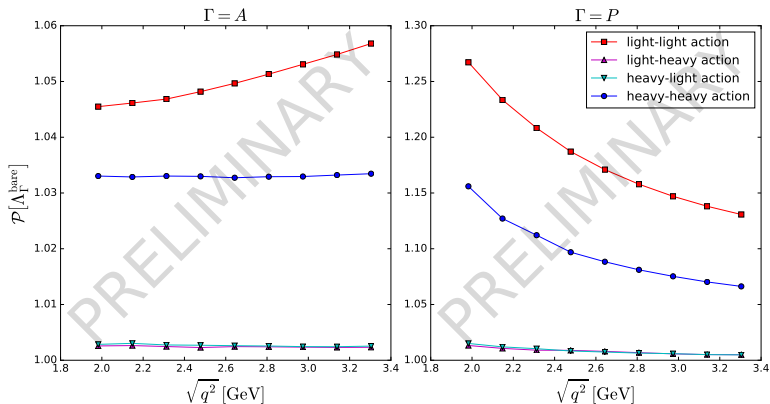
So for actions i, j

$$\frac{\mathcal{P}[\Lambda_A^{\text{bare}}]^{ii} \mathcal{P}[\Lambda_A^{\text{bare}}]^{jj}}{(\mathcal{P}[\Lambda_A^{\text{bare}}]^{ij})^2} = \frac{(Z_A^{ij})^2}{Z_A^{ii} Z_A^{jj}}$$

But for non-mixed actions we can determine Z_A^{ii} from conserved current.

Preliminary mixed action renormalisation

First study on single configuration

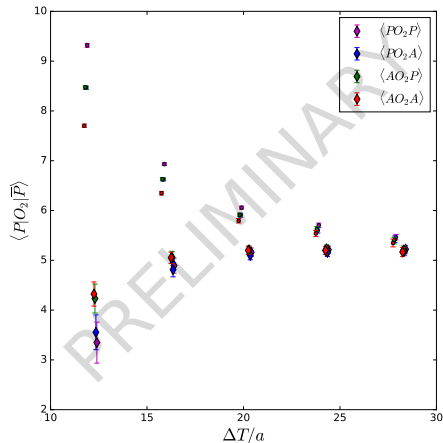


⇒ mixed NPR is feasible

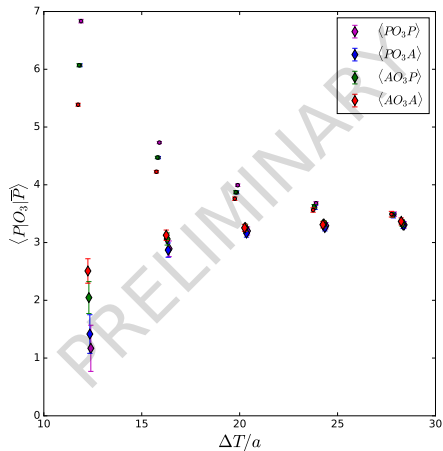
⇒ need to compute Z_A^{hh} from conserved current to obtain Z_A^{hl}

JLQCD + RBC/UKQCD data: additional operators

$VV - AA$

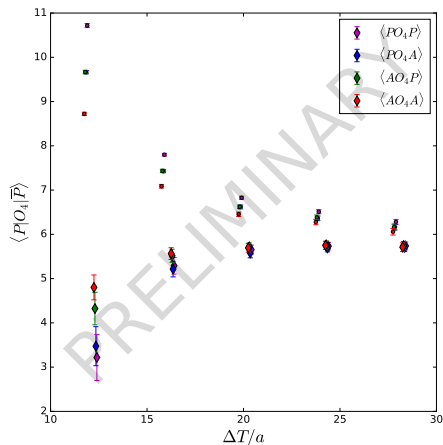


$SS + PP$



JLQCD + RBC/UKQCD data: additional operators

$SS - PP$



TT

