Form-factors for semi-leptonic B decays

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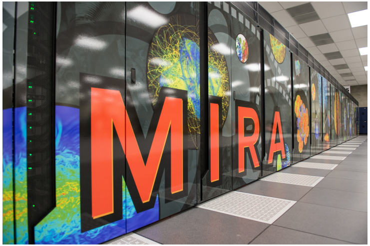
introduction

Nonperturbative Lattice QCD calculation

- Lattice QCD allows for first principle calculations
- Also in the nonperturbative regime
- Systematical procedures to improve uncertainties
- Requires large scale computing facilities



[ALCF]

Lattice QCD

- Wick-rotate to Eucledian time $t \to i\tau$
- Use path integral formalism

$$\langle \mathcal{O} \rangle_E = \frac{1}{Z} \int \mathcal{D}[\psi, \overline{\psi}]\, \mathcal{D}[U]\, \mathcal{O}[\psi, \overline{\psi}, U]\, e^{-S_E[\psi,\overline{\psi},U]}$$

- Discretize space-time

⇒ Large but finite dimensional path integral

→ Finite lattice spacing $a \to$ UV regulator

→ Finite volume of length $L \to$ IR regulator

[DiRAC]

Lattice QCD calculation

Simulate

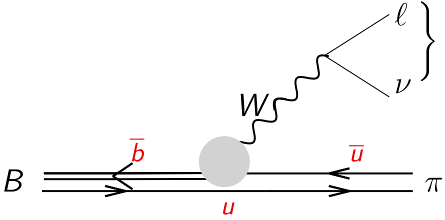
- at finite lattice spacing a
- in a finite volume L^3
  - $\Rightarrow$ discrete momenta $2\pi\vec{n}/L$
- lattice regularized
- bare input quark masses
  am_ℓ, am_s, am_c, am_b
  Mostly: $aM_\pi \neq aM_\pi^{\text{phys}}$

Desired result

- take $a \to 0$ limit
- take $L \to \infty$ limit
  - $\to$ continuous momenta $\vec{p}$
- match to some continuum scheme
- physical quark masses
  $m_l = m_{u/d}^{\text{phys}}$, $m_s = m_s^{\text{phys}}$, $m_c = m_c^{\text{phys}}$, $m_b = m_b^{\text{phys}}$

- Need to choose gauge and fermion action
- Need to control all limits keeping FV and discretization effects under control
  - $\to$ u quarks want large volume (large L^3) such that $M_\pi \cdot L > 4$
  - $\to$ b quarks want fine lattice (small a) i.e. $am_b \ll 1$

$|V_{ub}|$ from exclusive semileptonic $B \to \pi\ell\nu$ decay



$$q^2 = M_B^2 + M_\pi^2 - 2M_B E_\pi$$

- Conventionally parametrized by (B meson at rest)

$$\frac{d\Gamma(B \to \pi\ell\nu)}{dq^2} = \frac{G_F^2|V_{ub}|^2}{24\pi^3} \frac{(q^2 - m_\ell^2)^2\sqrt{E_\pi^2 - M_\pi^2}}{q^4 M_B^2}$$

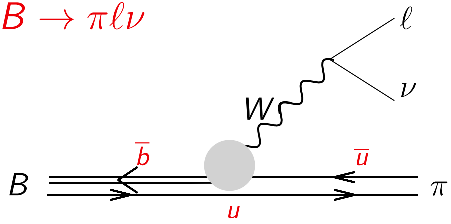
experiment CKM known

$$\times \left[\left(1 + \frac{m_\ell^2}{2q^2}\right) M_B^2(E_\pi^2 - M_\pi^2)|f_+(q^2)|^2 + \frac{3m_\ell^2}{8q^2}(M_B^2 - M_\pi^2)^2|f_0(q^2)|^2 \right],$$

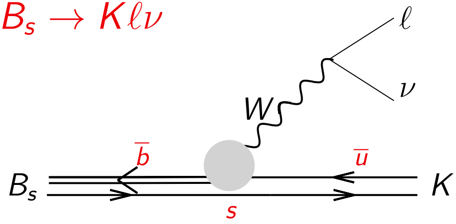
nonperturbative input

$|V_{ub}|$ from exclusive

$B \to \pi\ell\nu$



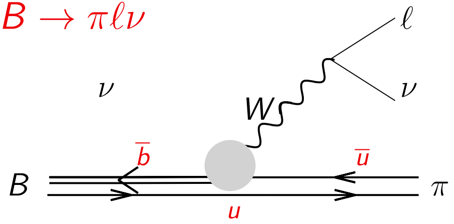
$B_s \to K\ell\nu$



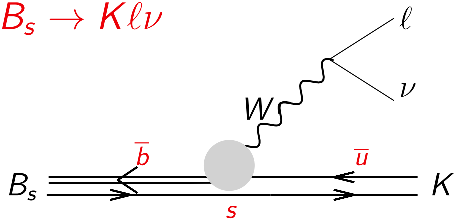
- Only spectator quark differs
- Lattice QCD prefers s quark over u quark: statistically more precise, computationally cheaper
- B factories run at $\Upsilon(4s)$ threshold $\Rightarrow$ B mesons
- LHC collisions create many B and B_s mesons which decay $\Rightarrow$ LHCb
  $\rightarrow$ LHCb prefers the ratio $(B_s \to D_s\ell\nu)/(B_s \to K\ell\nu) \Rightarrow |V_{cb}/V_{ub}|$

$|V_{ub}|$ from exclusive

$B \to \pi\ell\nu$

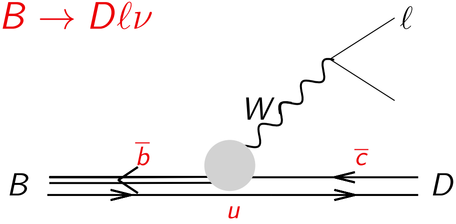


$B_s \to K\ell\nu$

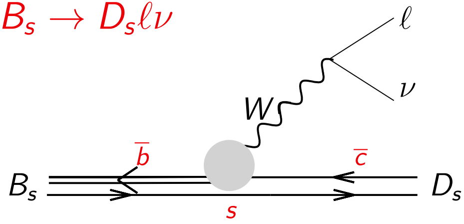


$|V_{cb}|$ from exclusive

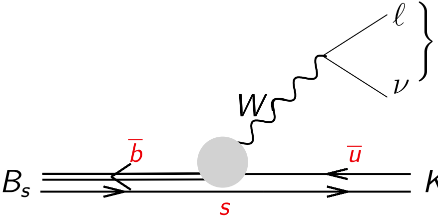
$B \to D\ell\nu$



$B_s \to D_s\ell\nu$



$|V_{ub}|$ from exclusive semileptonic $B_s \to K\ell\nu$ decay



$$q^2 = M_{B_s}^2 + M_K^2 - 2M_{B_s}E_K$$

- Conventionally parametrized by (B_s meson at rest)

$$\frac{d\Gamma(B_s \to K\ell\nu)}{dq^2} = \frac{G_F^2|V_{ub}|^2}{24\pi^3} \frac{(q^2 - m_\ell^2)^2\sqrt{E_K^2 - M_K^2}}{q^4 M_{B_s}^2}$$

experiment CKM known

$$\times \left[\left(1 + \frac{m_\ell^2}{2q^2}\right) M_{B_s}^2(E_K^2 - M_K^2)|f_+(q^2)|^2 + \frac{3m_\ell^2}{8q^2}(M_{B_s}^2 - M_K^2)^2|f_0(q^2)|^2\right]$$

nonperturbative input

Nonperturbative input

- Parametrizes interactions due to the (nonperturbative) strong force
- Use operator product expansion (OPE) to identify short distance contributions
- Calculate the flavor changing currents as point-like operators using lattice QCD

RBC-UKQCD's set-up

- RBC-UKQCD's 2+1 flavor domain-wall fermion and Iwasaki gauge action ensembles
  - → Three lattice spacings $a \sim 0.11$ fm, 0.08 fm, 0.07 fm; one ensemble with physical pions

    [PRD 78 (2008) 114509][PRD 83 (2011) 074508][PRD 93 (2016) 074505][JHEP 1712 (2017) 008]

RBC-UKQCD's set-up

	L	a^{-1}(GeV)	am_l	am_s	M_π(MeV)	# configs.	#sources	
C1	24	1.784	0.005	0.040	338	1636	1	[PRD 78 (2008) 114509]
C2	24	1.784	0.010	0.040	434	1419	1	[PRD 78 (2008) 114509]
M1	32	2.383	0.004	0.030	301	628	2	[PRD 83 (2011) 074508]
M2	32	2.383	0.006	0.030	362	889	2	[PRD 83 (2011) 074508]
M3	32	2.383	0.008	0.030	411	544	2	[PRD 83 (2011) 074508]
C0	*48*	*1.730*	*0.00078*	*0.0362*	*139*	*40*	*81/1*⋆	[PRD 93 (2016) 074505]
M0	64	2.359	0.000678	0.02661	139	—	—	[PRD 93 (2016) 074505]
F1	48	2.774	0.002144	0.02144	234	98	24	[JHEP 1712 (2017) 008]

⋆ All mode averaging: 81 "sloppy" and 1 "exact" solve [Blum et al. PRD 88 (2012) 094503]

- Lattice spacing determined from combined analysis [Blum et al. PRD 93 (2016) 074505]
- a: ~ 0.11 fm, ~ 0.08 fm, ~ 0.07 fm

RBC-UKQCD's set-up

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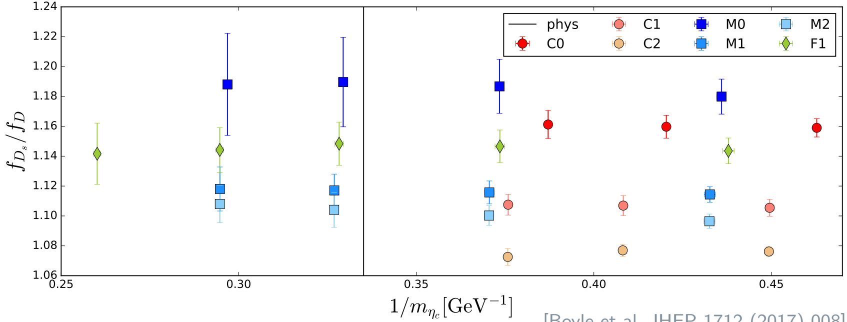
  - → Three lattice spacings $a \sim 0.11$ fm, 0.08 fm, 0.07 fm; one ensemble with physical pions
    [PRD 78 (2008) 114509][PRD 83 (2011) 074508][PRD 93 (2016) 074505][JHEP 1712 (2017) 008]
- Unitary and partially quenched domain-wall up/down quarks
  [Kaplan PLB 288 (1992) 342], [Shamir NPB 406 (1993) 90]
- Domain-wall strange quarks at/near the physical value
- Additional challenge $m_c = 1.28\text{GeV} \sim 270 \times m_d$

$$m_b = 4.18\text{GeV} \sim 1000 \times m_d$$

RBC-UKQCD's set-up

- Charm: Möbius domain-wall fermions optimized for heavy quarks [Boyle et al. JHEP 1604 (2016) 037]
  → Simulate 3 or 2 charm-like masses then extrapolate/interpolate



[Boyle et al. JHEP 1712 (2017) 008]

RBC-UKQCD's set-up

- Charm: Möbius domain-wall fermions optimized for heavy quarks [Boyle et al. JHEP 1604 (2016) 037]
  $\rightarrow$ Simulate 3 or 2 charm-like masses then extrapolate/interpolate

- Effective relativistic heavy quark (RHQ) action for bottom quarks
  [Christ et al. PRD 76 (2007) 074505], [Lin and Christ PRD 76 (2007) 074506]
  $\rightarrow$ Builds upon Fermilab approach [El-Khadra et al. PRD 55 (1997) 3933]
  $\rightarrow$ Allows to tune the three parameters ($m_0 a$, c_P, ζ) nonperturbatively [PRD 86 (2012) 116003]
  $\rightarrow$ Smooth continuum limit; heavy quark treated to all orders in $(m_b a)^n$
  $\rightarrow$ Mostly nonperturbative renormalization [Hashimoto et al. PRD61 (1999) 014502] [El-Khadra et al. PRD64 (2001) 014502]

$$Z_V^{bl} = \varrho\sqrt{Z_V^{ll} Z_V^{bb}}$$

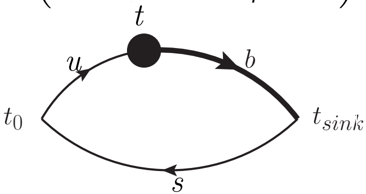
$B_s \to K\ell\nu$

in collaboration with
Jonathan M. Flynn, Ryan. C. Hill, Taku Izubuchi, Andreas Jüttner
Christoph Lehner, J. Tobias Tsang, Amarji Soni

$B_s \to K\ell\nu$ form factors

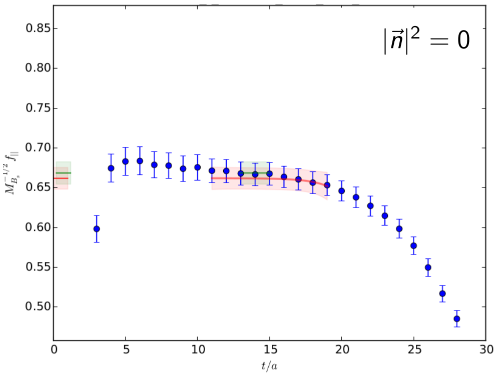
- Parametrize the hadronic matrix element for the flavor changing vector current V^μ in terms of the form factors $f_+(q^2)$ and $f_0(q^2)$

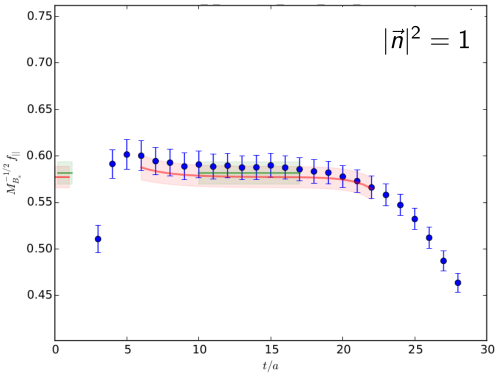
$$\langle K|V^\mu|B_s\rangle = f_+(q^2)\left(p^\mu_{B_s} + p^\mu_K - \frac{M^2_{B_s} - M^2_K}{q^2}q^\mu\right) + f_0(q^2)\frac{M^2_{B_s} - M^2_K}{q^2}q^\mu$$



- Calculate 3-point function by
  - → Inserting a quark source for a strange quark propagator at t_0
  - → Allow it to propagate to t_{sink}, turn it into a sequential source for a b quark
  - → Use a "light" quark propagating from t_0 and contract both at t with $t_0 \le t \le t_{sink}$

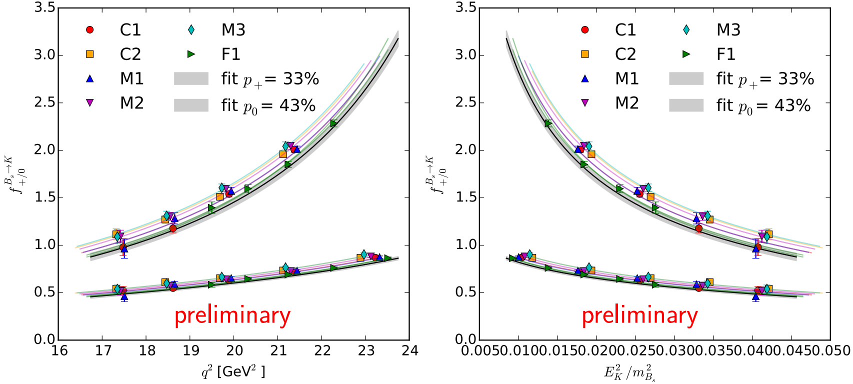
$B_s \to K\ell\nu$ form factors: F1 ensemble





- Comparison of fit to the ground state only with fit including one excited state term for K and B_s

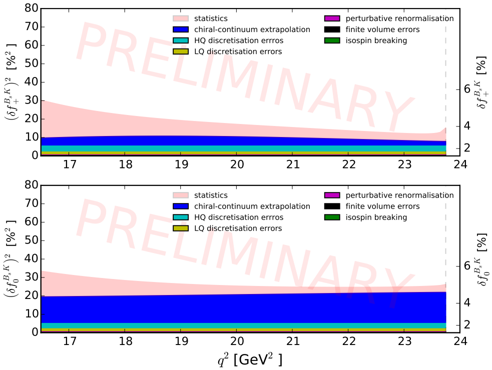
Chiral-continuum extrapolation using SU(2) hard-kaon χPT



- Updating calculation [PRD 91 (2015) 074510] with improved values for a^{-1} and RHQ parameters
- $f_{pole}(M_K, E_K, a^2) = \frac{1}{E_K+\Delta} c^{(1)} \times \left[1 + \frac{\delta f}{(4\pi f)^2} + c^{(2)} \frac{M_K^2}{\Lambda^2} + c^{(3)} \frac{E_K}{\Lambda} + c^{(4)} \frac{E_K^2}{\Lambda^2} + c^{(5)} \frac{a^2}{\Lambda^2 a_{32}^4}\right]$
- δf non-analytic logs of the kaon mass and hard-kaon limit is taken by $M_K/E_K \to 0$

Estimate systematic errors due to

- Chiral-continuum extrapolation
  - → Use alternative fit functions
  - → Impose different cuts on the data
- Discretization errors of light and heavy quarks
  - → Estimate via power-counting
- Uncertainty of the renormalization factors
  - → Estimate effect of higher loop corrections
- Finite volume, iso-spin breaking, ...
- Uncertainty due to RHQ parameters and lattice spacing (a^{-1})
  - → Carry out additional simulations to test effects on form factors
- Uncertainty of strange quark mass
  - → Repeat simulation with different valence quark mass



(plot still incomplete)

$\Longrightarrow$ full error budget

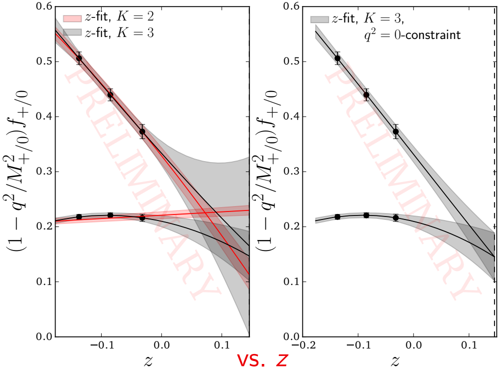
Kinematical extrapolation (z-expansion)

- Map q^2 to z with minimized magnitude in the semileptonic region: $|z| \leq 0.146$

$$z(q^2, t_0) = \frac{\sqrt{1-q^2/t_+} - \sqrt{1-t_0/t_+}}{\sqrt{1-q^2/t_+} + \sqrt{1-t_0/t_+}} \quad \text{with} \quad \begin{aligned} t_\pm &= (M_B \pm M_\pi)^2 \\ t_0 \equiv t_{\text{opt}} &= (M_B + M_\pi)(\sqrt{M_B} - \sqrt{M_\pi})^2 \end{aligned}$$

[Boyd, Grinstein, Lebed, PRL 74 (1995) 4603]

[Bourrely, Caprini, Lellouch, PRD 79 (2009) 013008]



vs. z

- Express f_+ as convergent power series
- f_0 is analytic, except for B^* pole
- BCL with poles $M_+ = B^* = 5.33$ GeV and $M_0 = 5.63$ GeV
- Exploit kinematic constraint $f_+ = f_0\big|_{q^2=0}$

→ Include HQ power counting to constrain size of f_+ coefficients

- Only some systematics included!

Kinematical extrapolation (z-expansion)

- Map q^2 to z with minimized magnitude in the semileptonic region: $|z| \leq 0.146$

$$z(q^2, t_0) = \frac{\sqrt{1-q^2/t_+} - \sqrt{1-t_0/t_+}}{\sqrt{1-q^2/t_+} + \sqrt{1-t_0/t_+}} \quad \text{with} \quad \begin{array}{l} t_\pm = (M_B \pm M_\pi)^2 \\ t_0 \equiv t_{opt} = (M_B + M_\pi)(\sqrt{M_B} - \sqrt{M_\pi})^2 \end{array}$$

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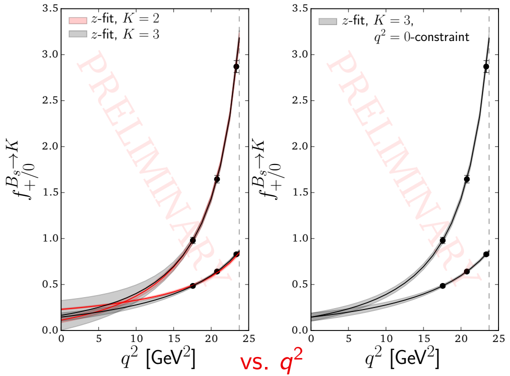
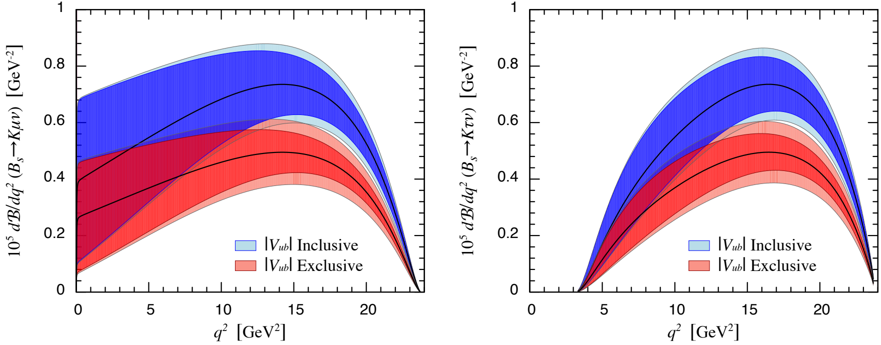


- Allows to compare shape of form factors
  - → Obtained by other lattice calculations
    [Bouchard et al. PRD 90 (2014) 054506]
    [Bazavov et al. arXiv:1901.02561]
  - → Predicted by QCD sum rules and alike
- Combination with experiment leads to the overall normalization: $|V_{ub}|$
- Only some systematics included!

Phenomenological interpretation (2015) [PRD 91 (2015) 074510]

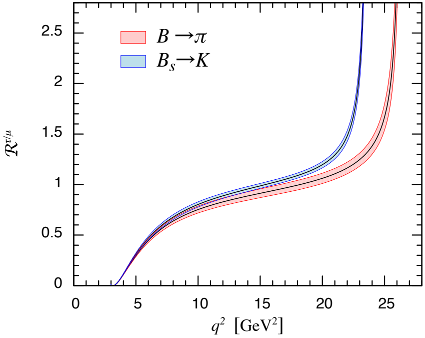
- Predict SM differential branching fractions using $|V_{ub}|$ as input for lepton $= \mu$ or τ



Phenomenological interpretation (2015) [PRD 91 (2015) 074510]

- Predict SM differential branching fractions using $|V_{ub}|$ as input for lepton $= \mu$ or τ

- Predict ratio of branching fractions $\rightsquigarrow$ LFUV

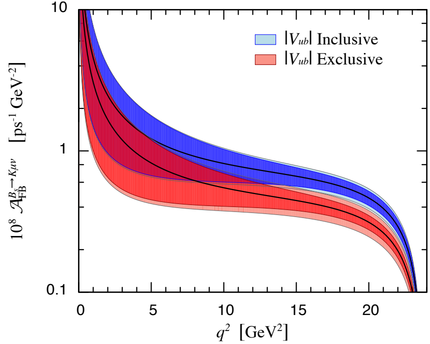


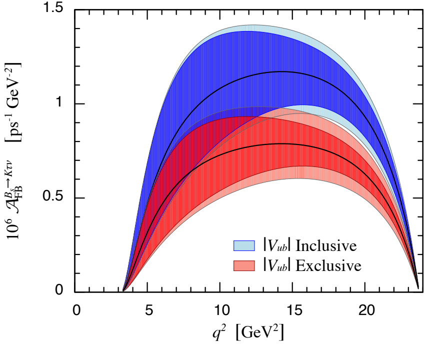
$$R_\pi^{\tau/\mu} = 0.69(19)$$

$$R_K^{\tau/\mu} = 0.77(12)$$

Phenomenological interpretation (2015) [PRD 91 (2015) 074510]

- Predict SM differential branching fractions using $|V_{ub}|$ as input for lepton $= \mu$ or τ
- Predict ratio of branching fractions $\rightsquigarrow$ LFUV
- Predict forward-backward asymmetries using $|V_{ub}|$ as input for lepton $= \mu$ or τ





$B_s \to D_s \ell\nu$

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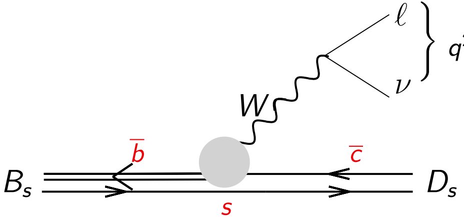


$$q^2 = M_{B_s}^2 + M_{D_s}^2 - 2M_{B_s}E_{D_s}$$

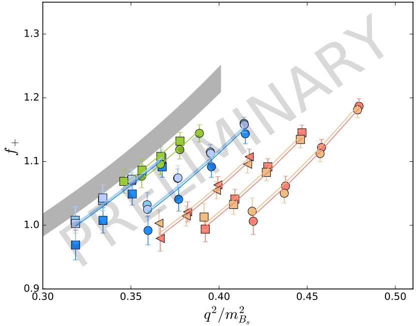
- Conventionally parametrized by (B_s meson at rest)

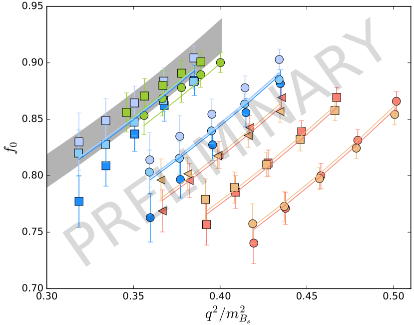
$$\underbrace{\frac{d\Gamma(B_s \to D_s\ell\nu)}{dq^2}}_{\text{experiment}} = \underbrace{\frac{G_F^2|V_{ub}|^2}{24\pi^3}}_{\text{CKM}} \underbrace{\frac{(q^2 - m_\ell^2)^2\sqrt{E_{D_s}^2 - M_{D_s}^2}}{q^4 M_{B_s}^2}}_{\text{known}}$$

$$\times \left[\left(1 + \frac{m_\ell^2}{2q^2}\right) M_{B_s}^2(E_{D_s}^2 - M_{D_s}^2)|f_+(q^2)|^2 + \frac{3m_\ell^2}{8q^2}(M_{B_s}^2 - M_{D_s}^2)^2|f_0(q^2)|^2\right]$$

nonperturbative input

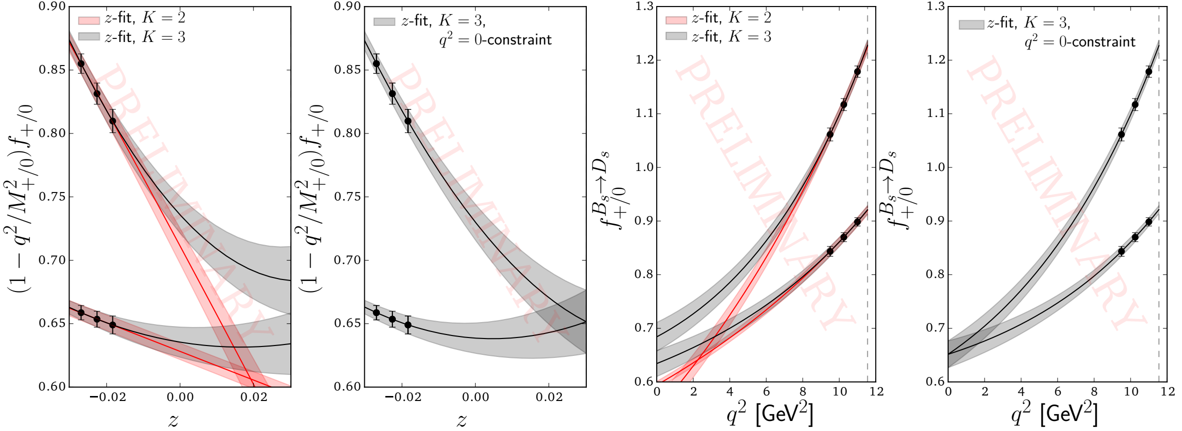
Global fit $B_s \to D_s\ell\nu$





- $f(q^2, a, M_\pi, M_{D_s}) = \dfrac{\alpha_0 + \alpha_1 M_{D_s} + \alpha_2 a^2 + \alpha_3 M_\pi^2}{1 + \alpha_4 q^2/M^2_{B_s}}$
- Extrapolation to the continuum limit with physical quark masses
- Error budget still work in progress

z-expansion



- BCL with poles $M_+ = B_c^* = 6.33$ GeV and $M_0 = 6.42$ GeV
- Comparison with CLN in progress

Status $B_s \to K\ell\nu$ and $B_s \to D_s\ell\nu$

- $B_s \to K\ell\nu$ chiral-continuum extrapolation
- $B_s \to D_s\ell\nu$ global fit (M_π, M_{D_s}, a^2, q^2)
- Extract synthetic data points

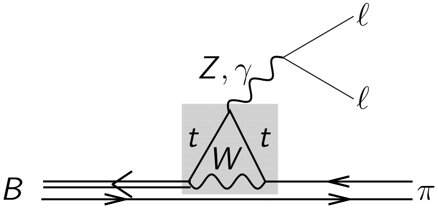
- Full systematic error budget
  - → RHQ parameter tuning
  - → Continuum extrapolation:
    cut to data set, different fit functions, ...
  - → Charm extrapolation
  - → FV, higher order disc. effects, isospin,
    s-quark mass tuning, ...

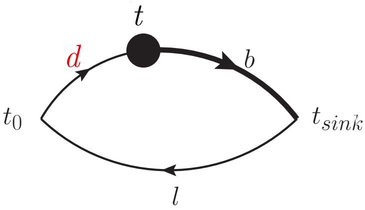
- z-expansion over full q^2 range
  - → BGL vs. BCL
  - → Test CLN for $B_s \to D_s\ell\nu$
  - → Number of synthetic data points
  - → Different truncations
  - → Incl. vs. excludig $f_+ = f_0\Big|_{q^2=0}$

- Phenomenology: $R(K)$, $R(D_s)$, ...

more form factors

$B \to \pi\ell^+\ell^-$ form factor (FCNC)





- If the daughter quark is a d-quark, we have a FCNC decay at loop-level
  $\to$ Need to implement additional operators
- Dominant contributions at short distance: f_0, f_+, and f_T

$$\langle\pi(k)|i\bar{d}\sigma^{\mu\nu}b(p)|B\rangle = 2\frac{p^\mu k^\nu - p^\nu k^\mu}{M_B+M_\pi} f_T(q^2)$$

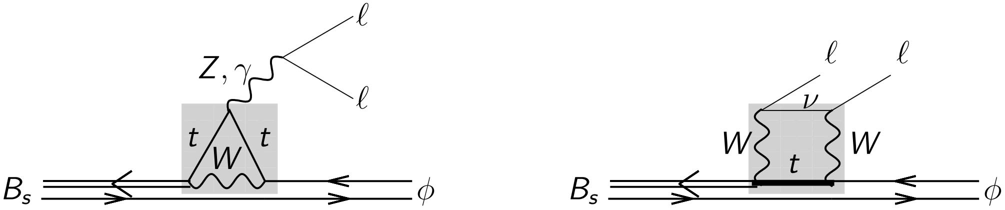
$B \to D^*\ell\nu$ form factors

- Vector final state with narrow width approximation

$$\langle D^*(k,\lambda)|\bar{c}\gamma^\mu b|B(p)\rangle = f_V \frac{2i\epsilon^{\mu\nu\rho\sigma}\varepsilon^*_\nu k_\rho p_\sigma}{M_B + M_D^*}$$

$$\begin{aligned}\langle D^*(k,\lambda)|\bar{c}\gamma^\mu\gamma_5 b|B(p)\rangle =& f_{A_0}(q^2)\frac{2M_{D^*}\varepsilon^*\cdot q}{q^2}q^\mu \\ &+ f_{A_1}(q^2)(M_B + M_{D^*})\left[\varepsilon^{*\mu} - \frac{\varepsilon^*\cdot q}{q^2}q^\mu\right] \\ &- f_{A_2}(q^2)\frac{\varepsilon^*\cdot q}{M_B + M_{D^*}}\left[k^\mu + p^\mu - \frac{M_B^2 - M_{D^*}^2}{q^2}q^\mu\right]\end{aligned}$$

$B_s \to \phi\ell^+\ell^-$ form factors (FCNC)



- Vector final state treated in narrow width approximation
- Effective Hamiltonian

$$\mathcal{H}_{\text{eff}}^{b\to s} = -\frac{4G_F}{\sqrt{2}} V_{ts}^* V_{tb} \sum_i^{10} C_i O_i^{(\prime)}$$

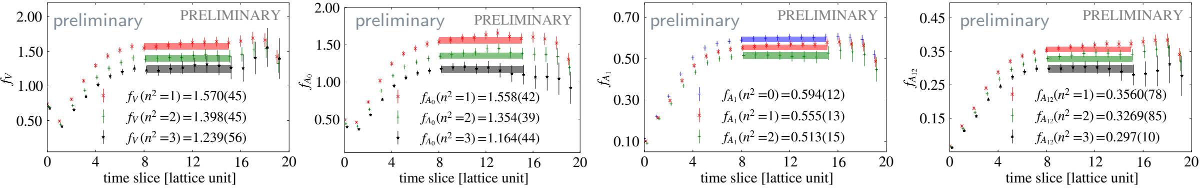
- Leading contributions at short distance

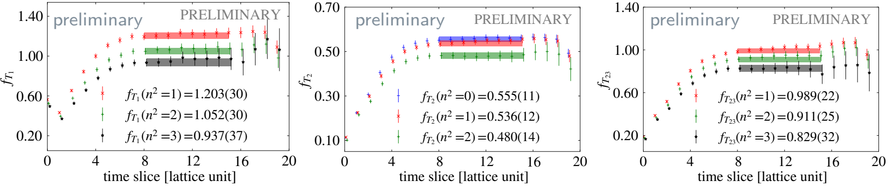
$$O_7^{(\prime)} = \frac{m_b e}{16\pi^2}\bar{s}\sigma^{\mu\nu}P_{R(L)}bF_{\mu\nu} \qquad O_9^{(\prime)} = \frac{e^2}{16\pi^2}\bar{s}\gamma^\mu P_{L(R)}b\bar{\ell}\gamma_\mu\ell \qquad O_{10}^{(\prime)} = \frac{e^2}{16\pi^2}\bar{s}\gamma^\mu P_{L(R)}b\bar{\ell}\gamma_\mu\gamma^5\ell$$

Seven form factors

$$\begin{aligned}
\langle\phi(k,\lambda)|\bar{s}\gamma^\mu b|B_s(p)\rangle =& f_V(q^2)\frac{2i\epsilon^{\mu\nu\rho\sigma}\varepsilon^*_\nu k_\rho p_\sigma}{M_{B_s}+M_\phi}\\
\langle\phi(k,\lambda)|\bar{s}\gamma^\mu\gamma_5 b|B_s(p)\rangle =& f_{A_0}(q^2)\frac{2M_\phi\varepsilon^*\cdot q}{q^2}q^\mu\\
&+f_{A_1}(q^2)(M_{B_s}+M_\phi)\left[\varepsilon^{*\mu}-\frac{\varepsilon^*\cdot q}{q^2}q^\mu\right]\\
&-f_{A_2}(q^2)\frac{\varepsilon^*\cdot q}{M_{B_s}+M_\phi}\left[k^\mu+p^\mu-\frac{M_{B_s}^2-M_\phi^2}{q^2}q^\mu\right]\\
q_\nu\langle\phi(k,\lambda)|\bar{s}\sigma^{\nu\mu}b|B_s(p)\rangle =& 2f_{T_1}(q^2)\epsilon^{\mu\rho\tau\sigma}\varepsilon^*_\rho k_\tau p_\sigma\ ,\\
q_\nu\langle\phi(k,\lambda)|\bar{s}\sigma^{\nu\mu}\gamma^5 b|B_s(p)\rangle =& if_{T_2}(q^2)\left[\varepsilon^{*\mu}(M_{B_s}^2-M_\phi^2)-(\varepsilon^*\cdot q)(p+k)^\mu\right]\\
&+if_{T_3}(q^2)(\varepsilon^*\cdot q)\left[q^\mu-\frac{q^2}{M_{B_s}^2-M_\phi^2}(p+k)^\mu\right]
\end{aligned}$$

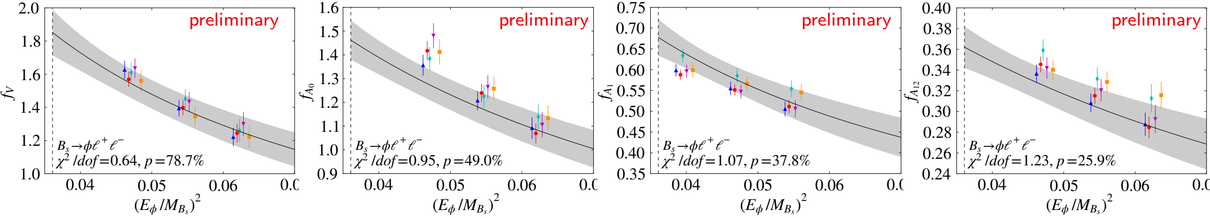
$B_s \to \phi\ell\ell$: seven form factors ($a^{-1} = 1.784$ GeV, $am_l^{sea} = 0.005$, $am_s = 0.03224$)

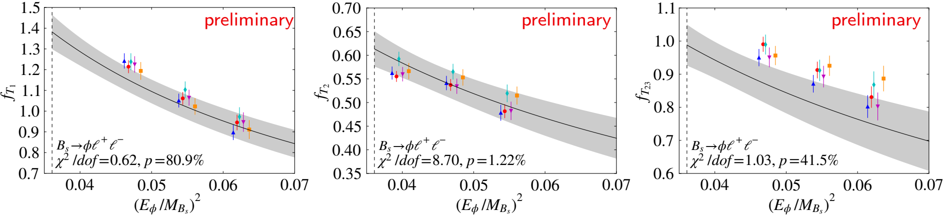




plots:
Edwin Lizarazo

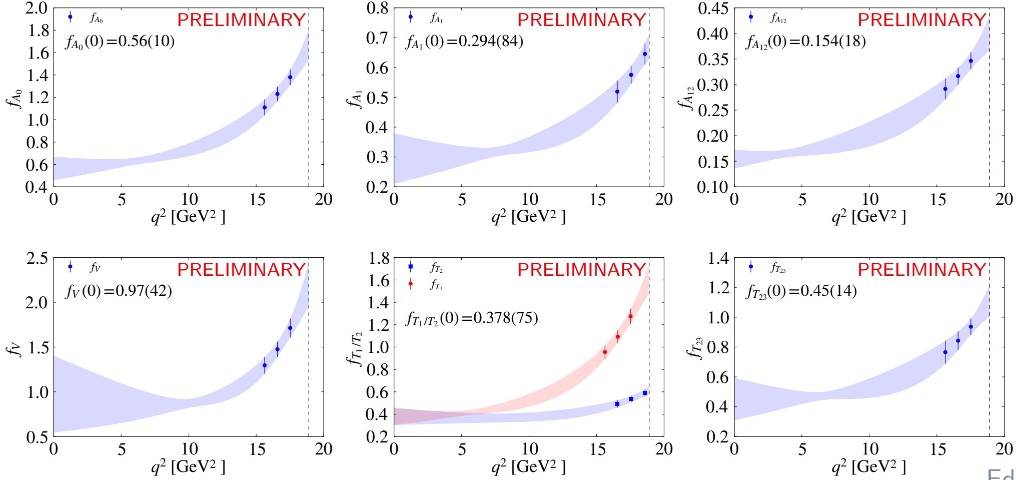
$B_s \to \phi\ell^+\ell^-$: seven form factors vs. q^2





plots:
Edwin Lizarazo

$B_s \to \phi\ell^+\ell^-$: first attempt to use z-parametrization



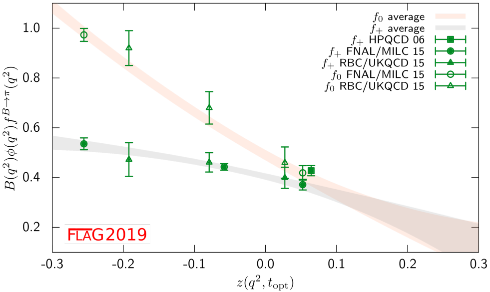
plots: Edwin Lizarazo

Flavor Lattice Averaging Group

[FLAG 2019]

$B \to \pi\ell\nu$

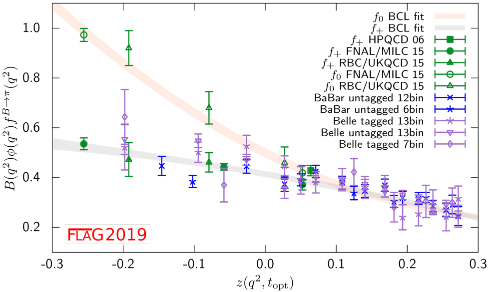
Collaboration	Ref.	N_f	publication status	continuum extrapolation	chiral extrapolation	finite volume	renormalization	heavy-quark treatment	z-parameterization	$\Delta\zeta^{B\pi}$
FNAL/MILC 15	[575]	2+1	A	★	○	★	○	✓	BCL	n/a
RBC/UKQCD 15	[576]	2+1	A	○	○	○	○	✓	BCL	1.77(34)
HPQCD 06	[573]	2+1	A	○	○	○	○	✓	n/a	2.07(41)(39)

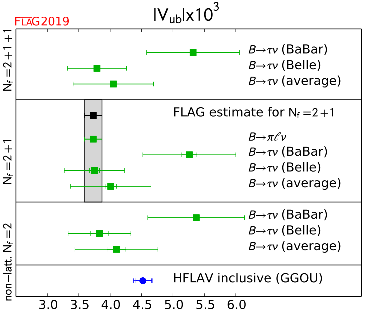


- Summary table indicating quality of key features
- Combined analysis accounting for correlations (if needed)
- FNAL/MILC also $B \to \pi\ell^+\ell^-$ (FCNC)
- Please do cite calculations feeding into FLAG averages

[FLAG 2019]

$B \to \pi\ell\nu$

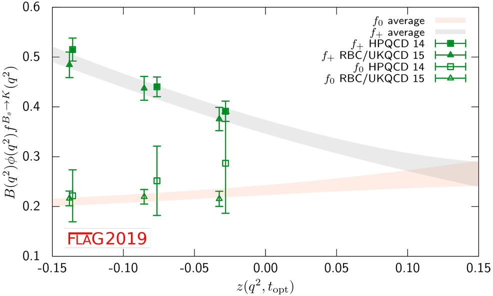


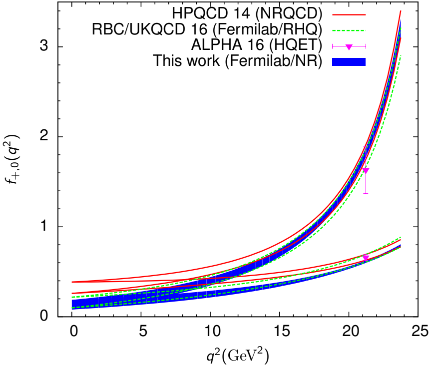


- Combination of lattice average with experimental results
- Extraction of $|V_{ub}|$
- FNAL/MILC also $B \to \pi\ell^+\ell^-$ (FCNC)
- Please do cite calculations feeding into FLAG averages

[FLAG 2019]

$B_s \to K\ell\nu$



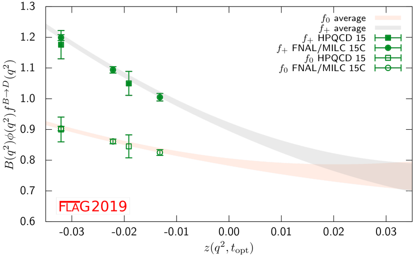


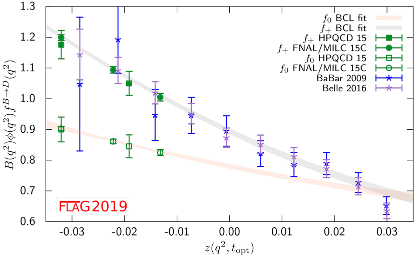
- New FNAL/MILC [arXiv:1901.02561]

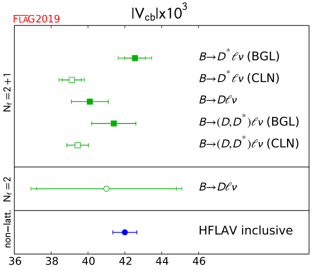
- Please do cite calculations feeding into FLAG averages

[FLAG 2019]

$B \to D\ell\nu$







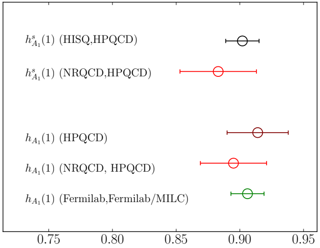
- Please do cite calculations feeding into FLAG averages

[FLAG 2019]

$B \to D_{(s)}^{(*)}\ell\nu$

Collaboration	Ref.	N_f	publication status	continuum extrapolation	chiral extrapolation	finite volume	renormalization	heavy-quark treatment	$w=1$ form factor / ratio	
HPQCD 15, HPQCD 17	[614, 616]	2+1	A	○	○	○	○	✓	$\mathcal{G}^{B\to D}(1)$	1.035(40)
FNAL/MILC 15C	[613]	2+1	A	★	○	★	○	✓	$\mathcal{G}^{B\to D}(1)$	1.054(4)(8)
Atoui 13	[610]	2	A	★	○	★	—	✓	$\mathcal{G}^{B\to D}(1)$	1.033(95)
HPQCD 15, HPQCD 17	[614, 616]	2+1	A	○	○	○	○	✓	$\mathcal{G}^{B_s\to D_s}(1)$	1.068(40)
Atoui 13	[610]	2	A	★	○	★	—	✓	$\mathcal{G}^{B_s\to D_s}(1)$	1.052(46)
HPQCD 17B	[618]	2+1+1	A	○	★	★	○	✓	$\mathcal{F}^{B\to D^*}(1)$	0.895(10)(24)
FNAL/MILC 14	[612]	2+1	A	★	○	★	○	✓	$\mathcal{F}^{B\to D^*}(1)$	0.906(4)(12)
HPQCD 17B	[618]	2+1+1	A	○	★	★	○	✓	$\mathcal{F}^{B_s\to D_s^*}(1)$	0.883(12)(28)
HPQCD 15, HPQCD 17	[614, 616]	2+1	A	○	○	○	○	✓	$R(D)$	0.300(8)
FNAL/MILC 15C	[613]	2+1	A	★	○	★	○	✓	$R(D)$	0.299(11)

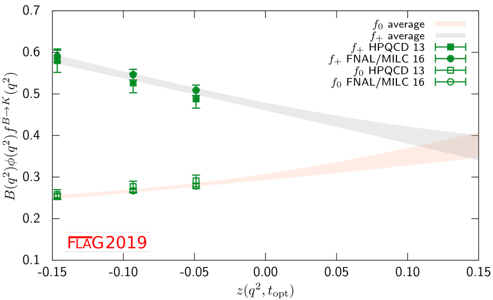
- New HPQCD $B_s \to D_s^*\ell\nu$ [arXiv:1904.02046]

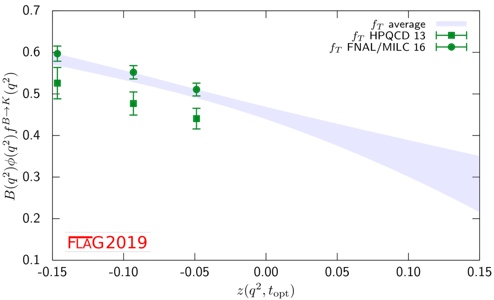


- Please do cite calculations feeding into FLAG averages

[FLAG 2019]

$B \to K\ell^+\ell^-$ (FCNC)





- Please do cite calculations feeding into FLAG averages

[FLAG 2019]

Other calculations (likely incomplete)

- FCNC decays with vector final state
  - → Seven form factors $B \to K^*\ell^+\ell^-$ and $B_s \to \phi\ell^+\ell^-$

    [Horgan, Liu, Meinel, and Wingate PRD89 (2014) 090501][PoS Lattice2014 (2015) 372]
  - → Angular analysis [PRL112 (2014) 212003]

- Exclusive baryonic decays
  - → $\Lambda_b \to \Lambda_c\ell\nu$ and $\Lambda_b \to p\ell\nu \Rightarrow |V_{cb}|/|V_{ub}|$ [Detmold, Lehner, Meinel, PRD92 (2015) 034503]
  - → $\Lambda_b \to \Lambda_c\tau\nu$ [Datta et al. JHEP08(2017)131]

outlook

Outlook

- Many calculation for exclusive decays are on it's way
- Updates will be reported Lattice 2019 (June 16–22, 2019)

- New ideas to compute inclusive decays using lattice techniques

  [Hashimoto PTEP 2017 (2017) 053B03] [Hansen, Meyer, Robaina arXiv:1704.08993]

Outlook: RBC-UKQCD

- Second (third) entirely independent analysis completed
- In the final stages to complete $B_s \to K\ell\nu$ and $B_s \to D_s\ell\nu$ form factor calculation
  - → As usual, carefully estimating all systematic uncertainties is tedious
- Our lattice calculation also includes
  - → $B \to \pi\ell\nu$, $B \to \pi\ell^+\ell^-$
  - → $B \to K^*\ell^+\ell^-$
  - → $B \to D^{(*)}\ell\nu$
  - → $B_s \to K^*\ell^+\ell^-$
  - → $B_s \to D_s^*\ell\nu$
  - → $B_s \to \phi\ell^+\ell^-$
  - → ...
- Current status $B_s \to K\ell\nu$ and $B_s \to D_s\ell\nu$: [arXiv:1903.02100]
- Future
  - → Add $48^3 \times 96$ ensemble with physical pions
- Parallel efforts: SU(3) breaking ratios [arXiv:1812.08791]

Resources for RBC-UKQCD's calculation

USQCD: Ds, Bc, and pi0 cluster (Fermilab), qcd12s cluster (Jlab), skylake cluster (BNL)

RBC qcdcl (RIKEN) and cuth (Columbia U)

UK: ARCHER, Cirrus (EPCC) and DiRAC (UKQCD)